

Pairing by nematic fluctuations in doped FeSe

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NPJ Quantum Materials 9, 28 (2024)
PRB 111, 094125 (2025)
arXiv:2605.05433 (2026)

Kazi Islam, UMN -> UW-Milwaukee

The story will be about
superconductivity

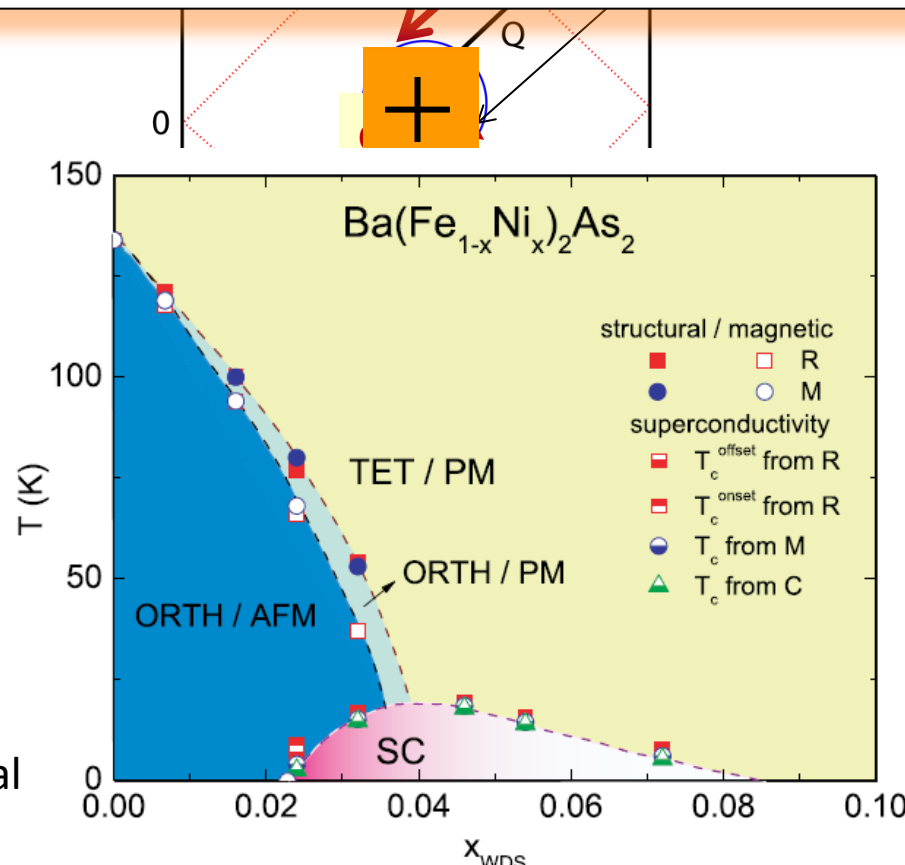
due to electron-electron interaction near
the onset of a nematic order in doped FeSe

The most discussed scenario of electronic pairing in correlated electron systems is spin fluctuation mechanism



Workhorse for SC in Fe-based materials from 2008

One needs strong (π, π) spin fluctuations to enhance pair hopping



Mazin et al, ...

Both inter
exceeds i

Canfield et al

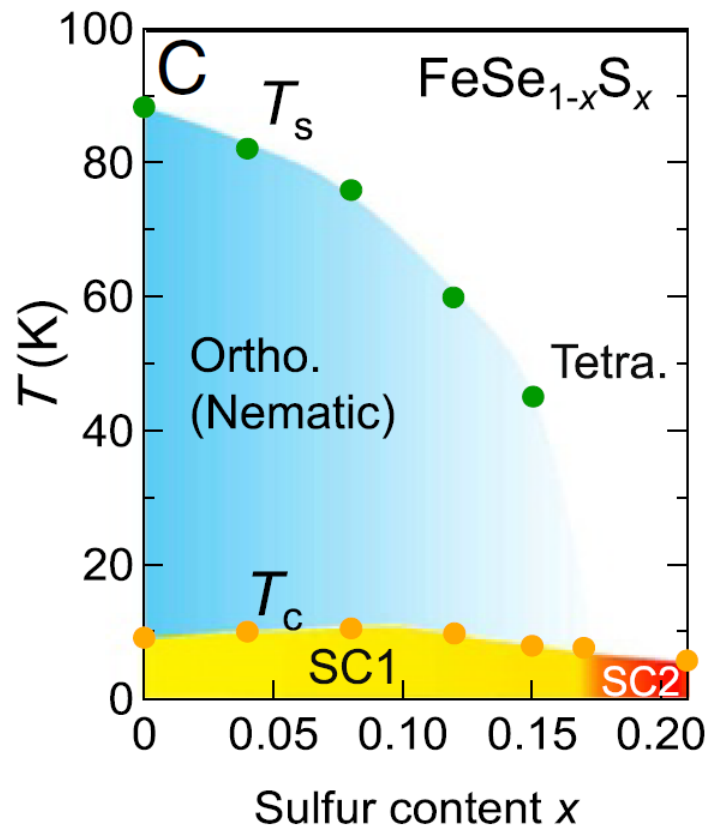
ket g_2 (pair hopping)

otliar, Nature 601, 35-44 (2022)

hopping g_2
y emerges
in.

Iron chalcogenide: S and Te-doped FeSe

$\text{FeSe}_{1-x}\text{S}_x$ and $\text{FeSe}_{1-x}\text{Te}_x$

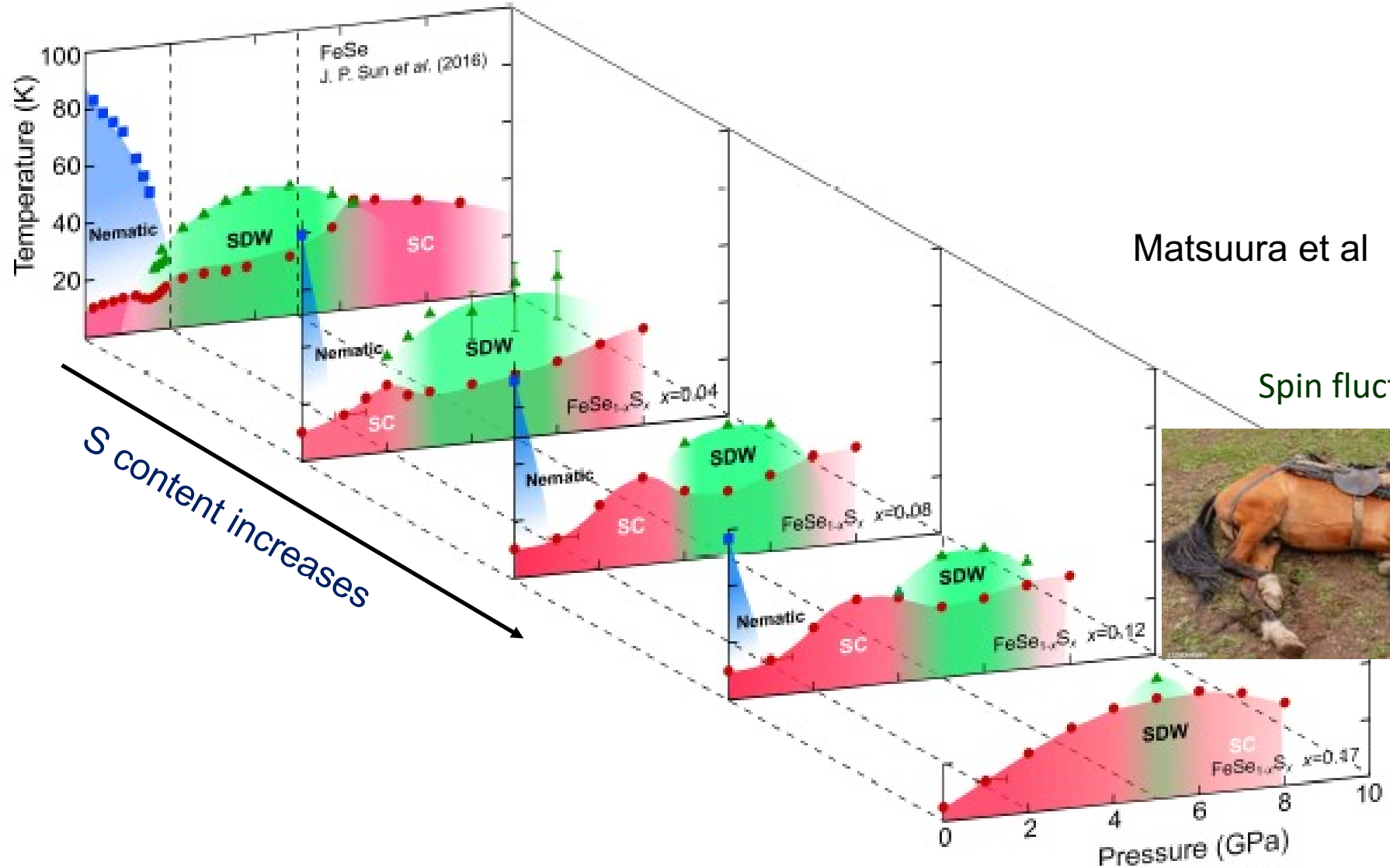


- ➔ Pure FeSe develops a nematic order at $T_n=85\text{K}$, SC at 9K
- ➔ Upon doping by S, T_n decreases and vanishes at around $x_c=0.18$
- ➔ At a critical x_c , superconducting T_c is finite
- ➔ No SDW at ambient pressure

$\text{FeSe}_{1-x}\text{S}_x$: magnetic fluctuations are reasonably strong at $x=0$, but are much weaker at $x=0.18$, where nematic order ends



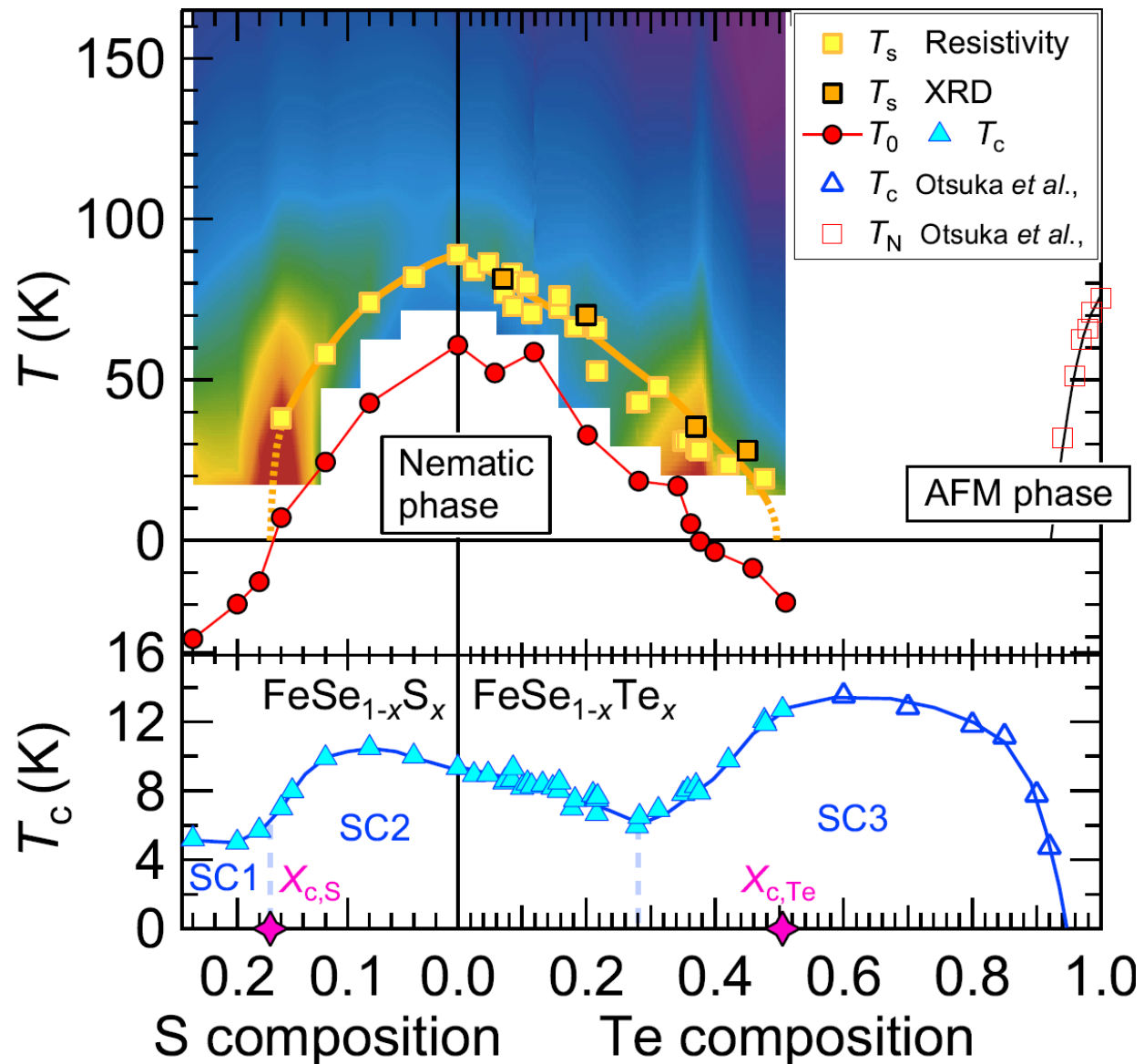
Spin fluctuations



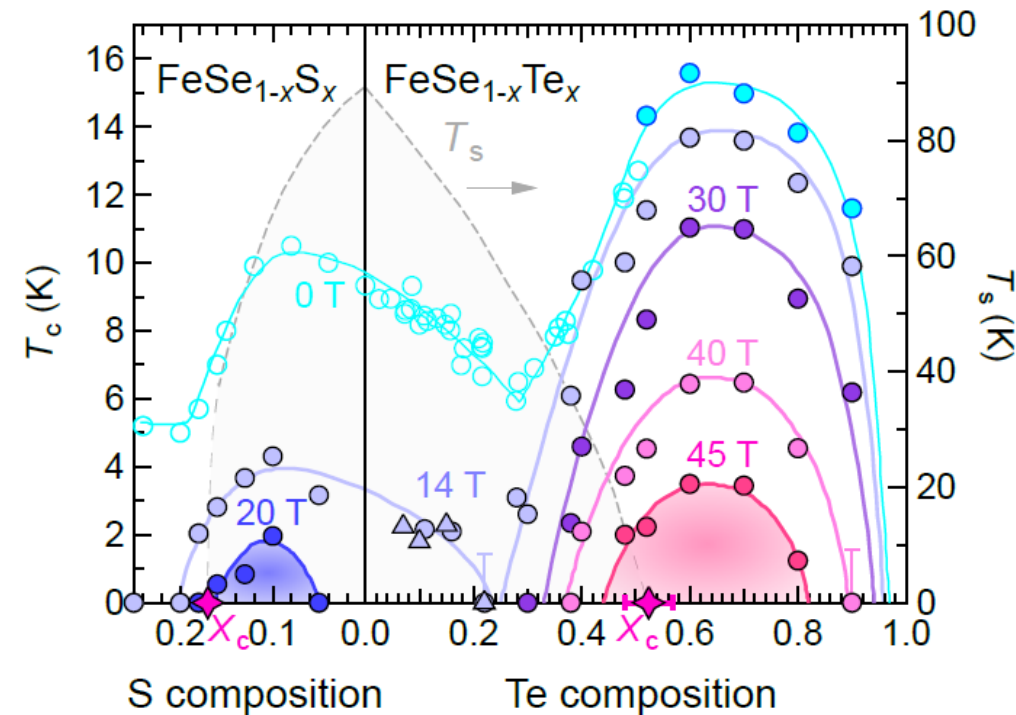
Spin fluctuations



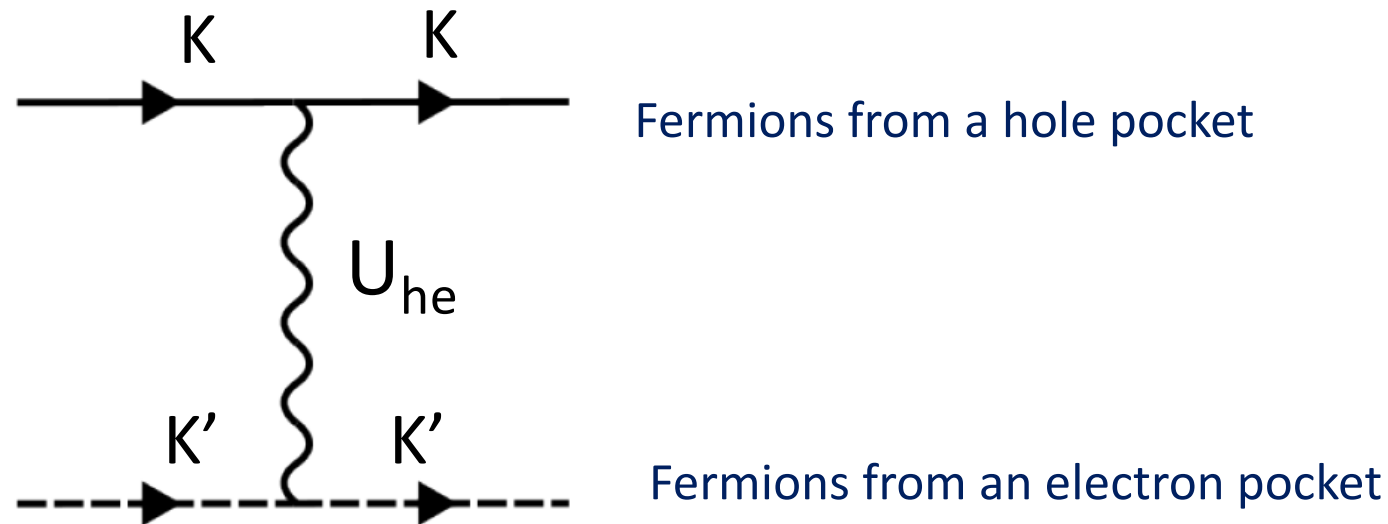
FeSeS and FeSeTe



Nematic fluctuations



The interaction that gives rise to a d-wave nematic order
and to SC,
mediated by nematic fluctuations,
is inter-pocket density-density repulsion



- This interaction is with zero momentum transfer and the total momentum near (π, π) .

Theory kitchen



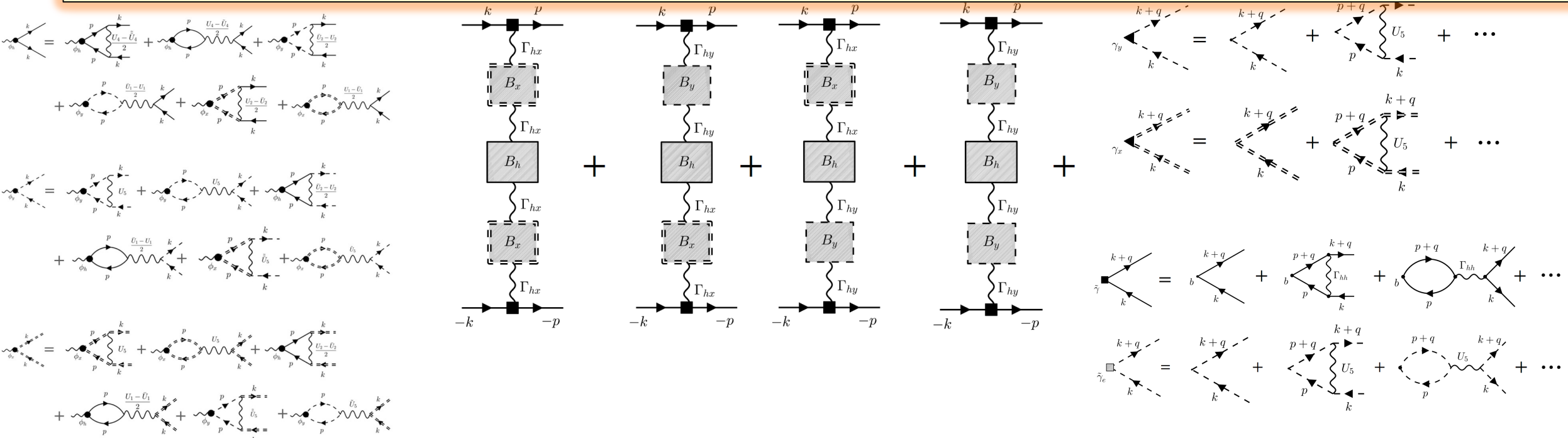
Pairing interaction within a given pocket

$$B_h = \text{loop with } \ell \text{ and } \ell+k-p \text{ and } \Gamma_{hh} \text{ interaction} + \dots$$

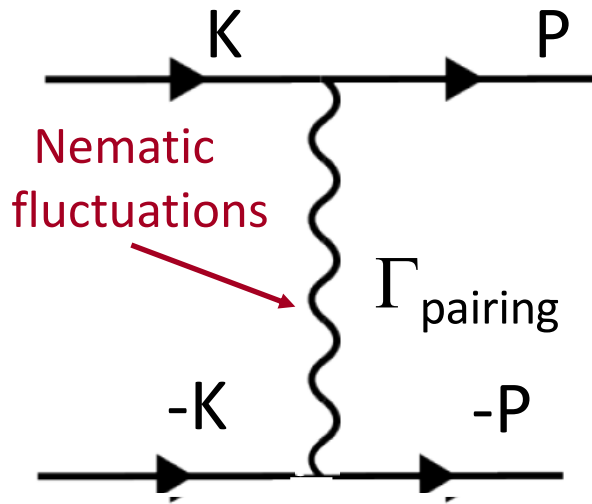
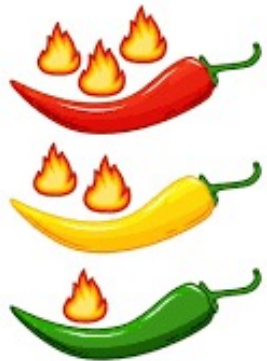
$$B_y = \text{loop with } \ell \text{ and } \ell+k-p \text{ and } U_5 \text{ interaction} + \dots$$

Inter-pocket density-density interaction U_{he} gives rise to:

- (i) d-wave nematic order changing sign between hole and electron pockets, and
- (ii) attractive pairing interaction, proportional to the d-wave χ_{nem}



The spicy part



Pairing by soft nematic fluctuations has been analyzed phenomenologically by several groups

Schattner, Lederer, Kivelson, Berg, Kontani, Onari, Schmalian, Garst, Fernandes, Klein, A.C.....

Phenomenological theories emphasized that the pairing interaction between fermions on the Fermi surface, specified by angles θ_k (for K) and θ_p (for P), is a function of $\theta_k - \theta_p$.
Different pairing channels (s,p,d...) are almost degenerate

Microscopic calculation: the attractive interaction, induced by U_{ph} , knows about the orbital content of low-energy excitations on the hole pocket (dxz and dyz)

Interaction is dressed by coherence factors of the transformation from orbitals to bands

$$\Gamma(\theta_k, \theta_p) = \Gamma(\theta_k - \theta_p, \theta_k + \theta_p) = - \underbrace{\cos^2(\theta_k + \theta_p)}_{\text{from coherent factors}} f(\theta_k - \theta_p)$$

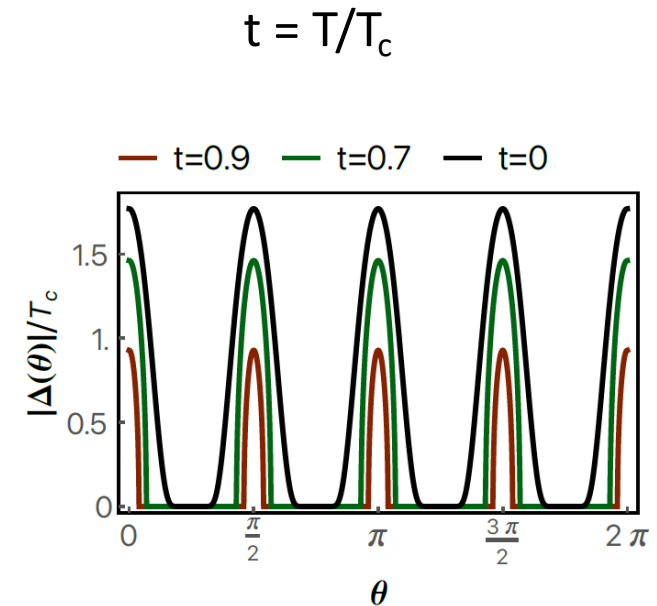
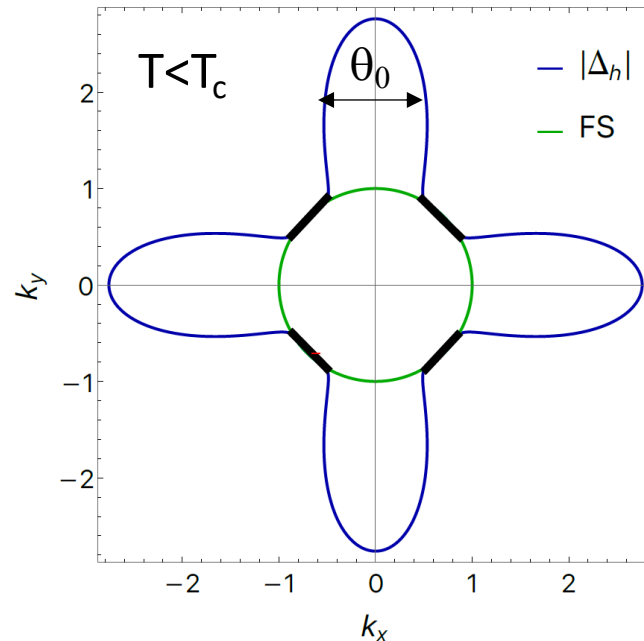
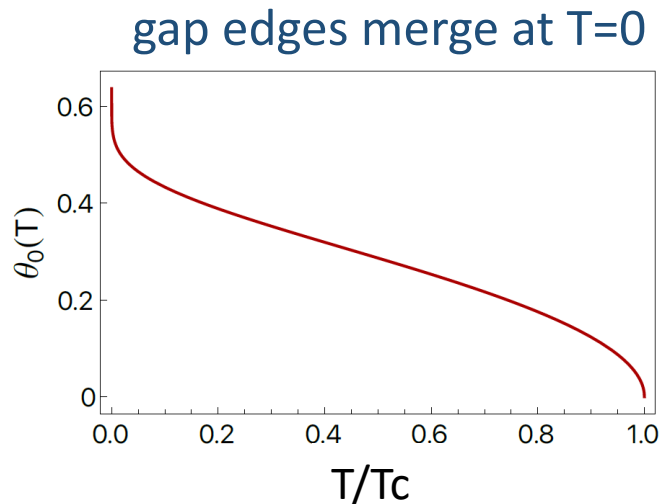
from coherent factors

Near a nematic QCP, relevant transferred momenta are small
(a nematic order has $Q=0$), and the pairing interaction effectively becomes

$$\Gamma(\theta_k, \theta_p) \approx \Gamma(\theta_k) = \underbrace{-g \cos^2(2\theta_k)}_{\text{from coherent factors}} \delta(\theta_k - \theta_p)$$

Interaction strength depends on the direction of momentum along the hole Fermi surface

SC gap opens at T_c only at $k=0, \pi/2, \dots$,
and extends along the FS as T decreases



Experiment

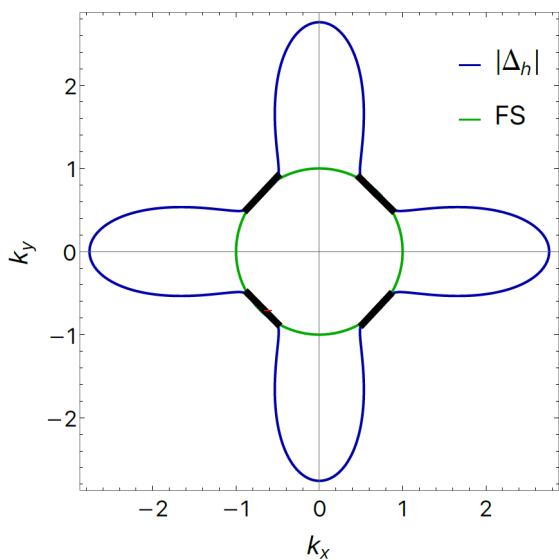




Can experiments distinguish between nematic and spin fluctuation scenarios?

Spin fluctuation and nematic pairing mechanisms yield qualitatively different gap anisotropy

Nematic scenario: a highly anisotropic gap function, with deep minima at 45° with respect to x and y directions in 1Fe unit cell

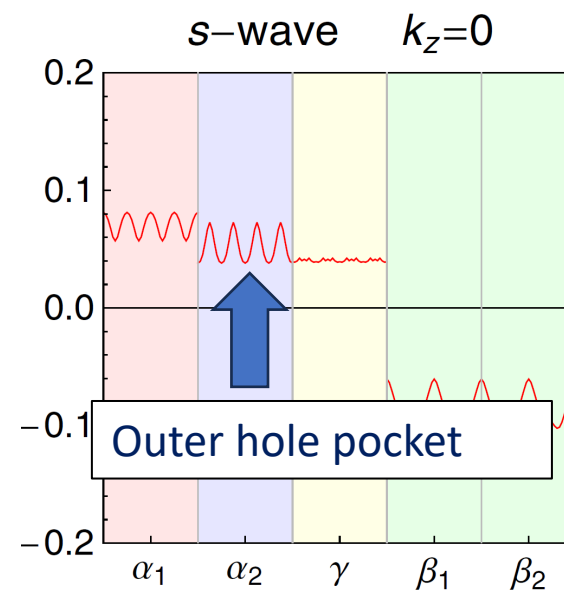


$c > 0$, gap minima along diagonals

Magnetic scenario: gap minima along x and y directions in 1Fe unit cell

Rotation of the gap min/max by 45°

$$\Delta = a + c \cos(4\Theta)$$

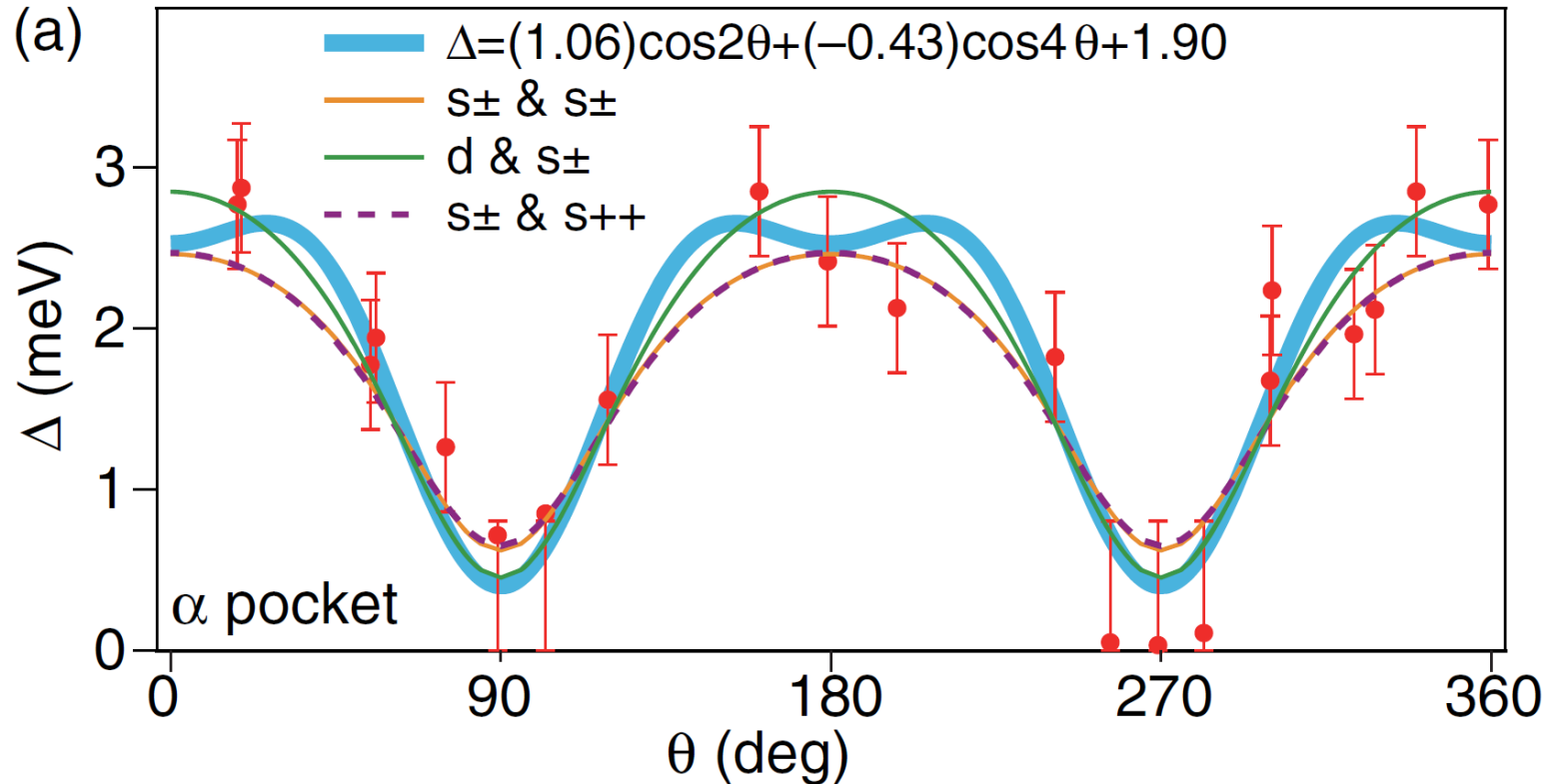


Graser et al

$c < 0$, gap minima along x and y

Gap anisotropy in the nematic state of $\text{FeSe}_{1-x}\text{S}_x$

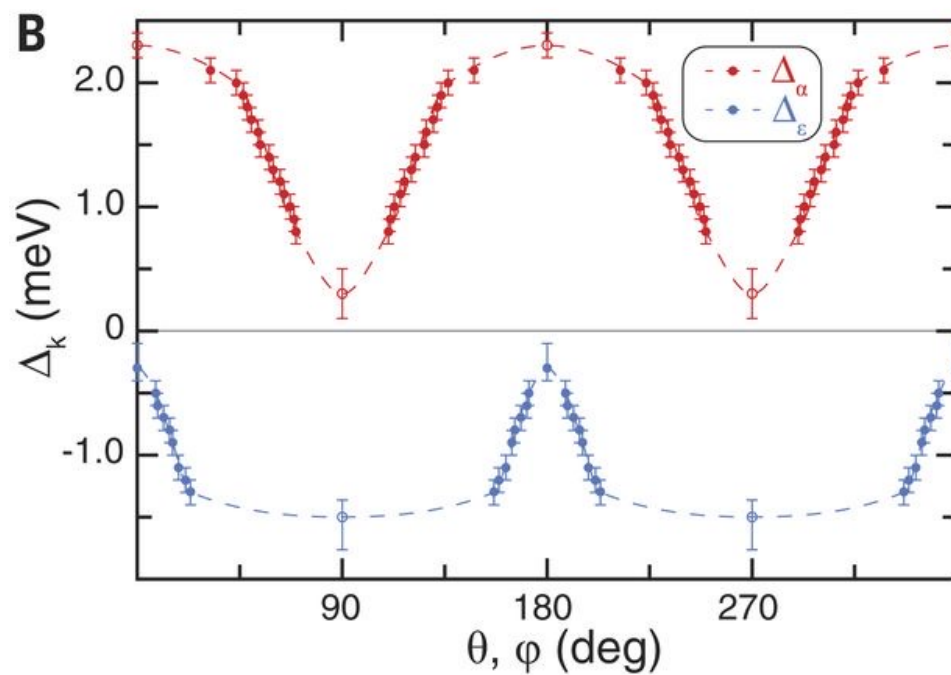
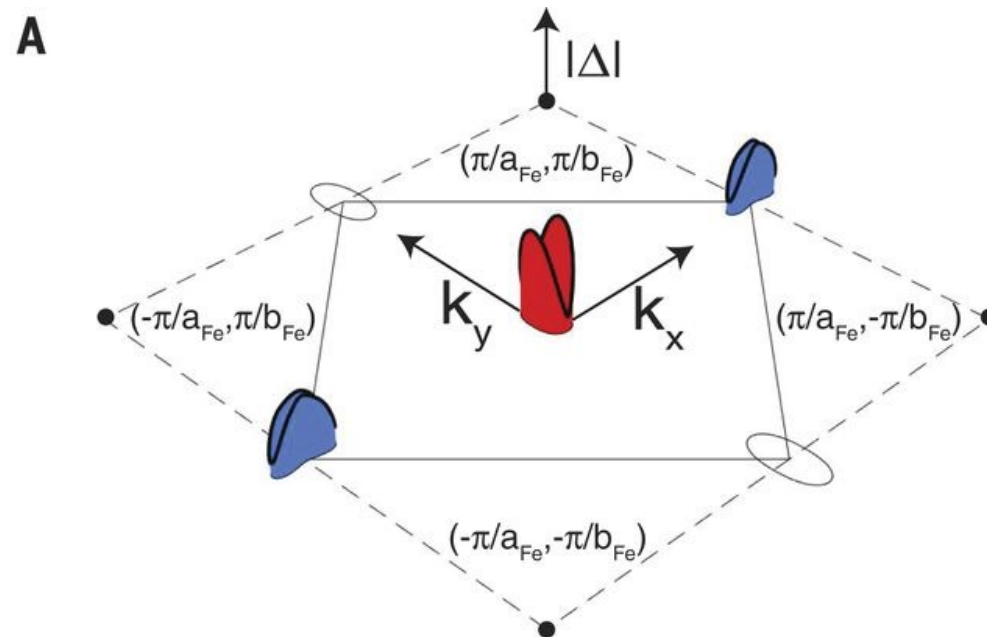
$$\Delta = a + b \cos(2\Theta) + c \cos(4\Theta)$$



D. Feng et al

Similar results:
X. J. Zhou et al

c is negative, spin fluctuations in the driver's seat



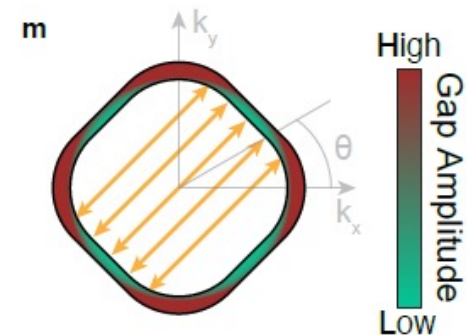
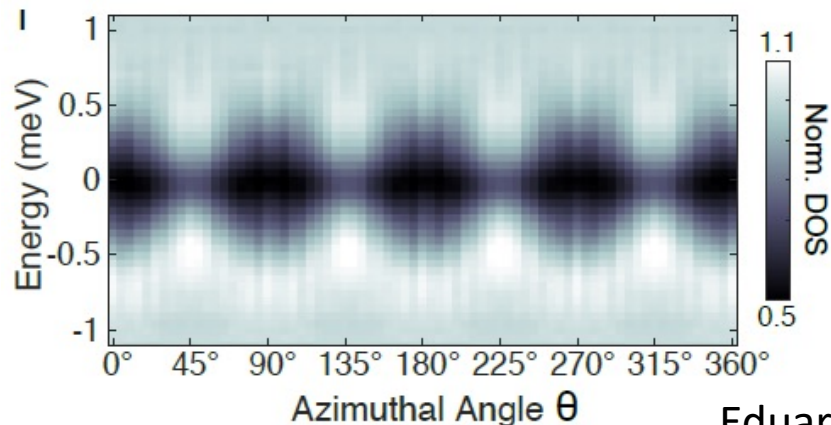
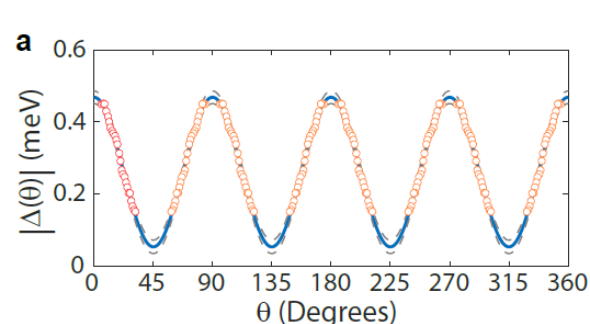
STM

Sprau et al

The same conclusion

Gap anisotropy in the tetragonal state (x = 0.19)

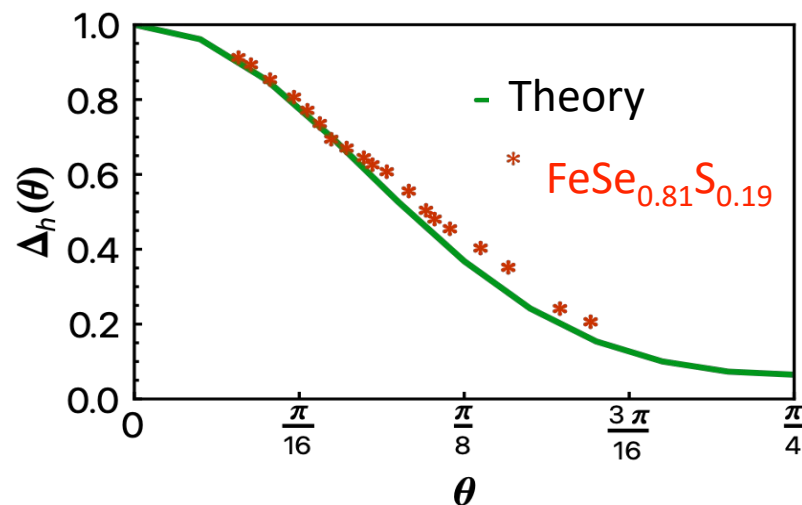
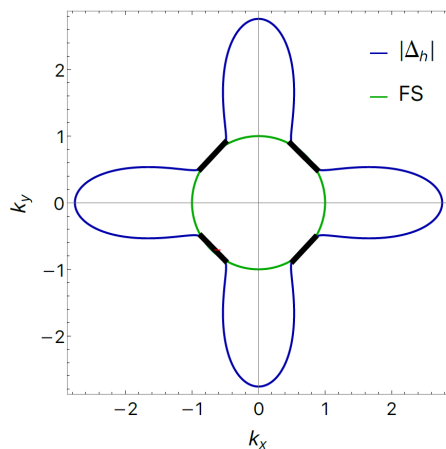
Gap anisotropy in $\text{FeSe}_{0.81}\text{S}_{0.19}$ from QPI



Eduardo H da Silva Neto, R. Fernandes....

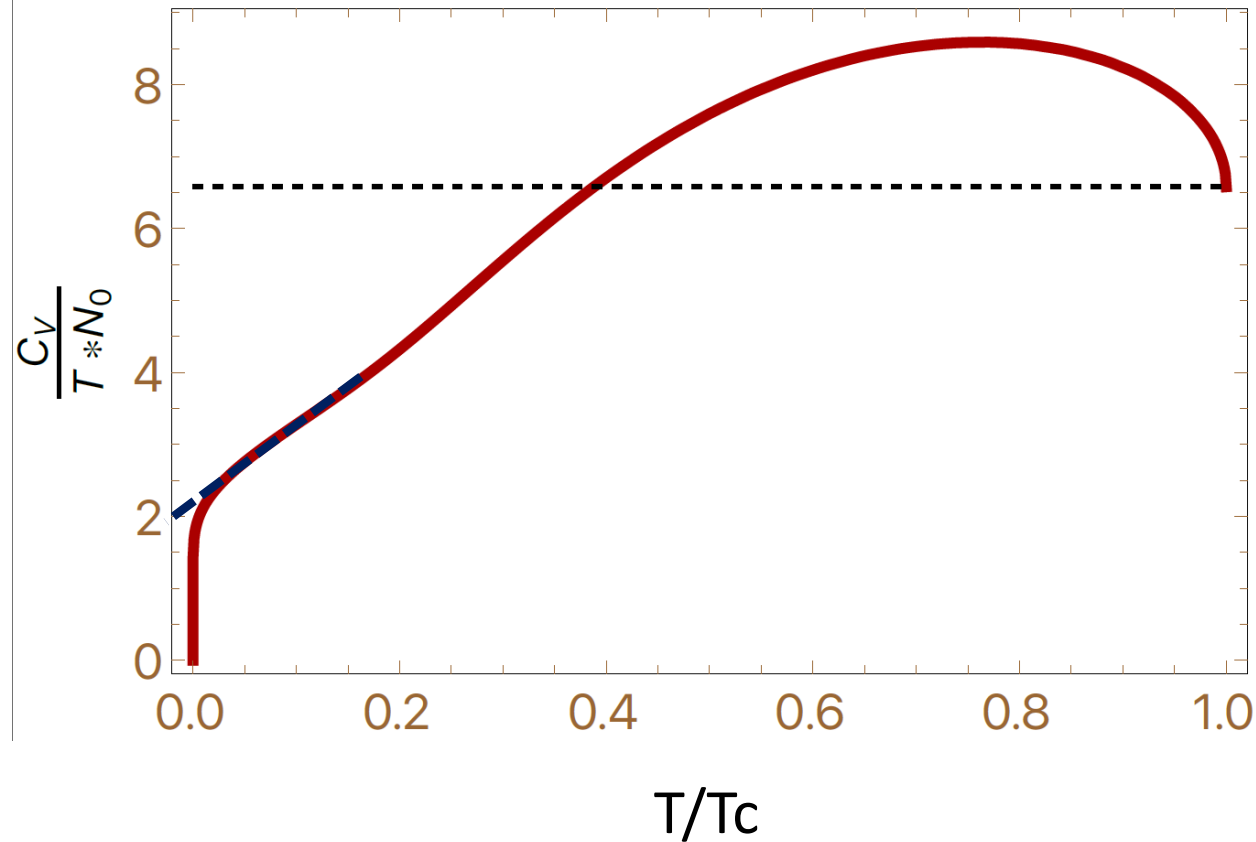
$$\Delta = a + c \cos(4\Theta)$$

c is positive

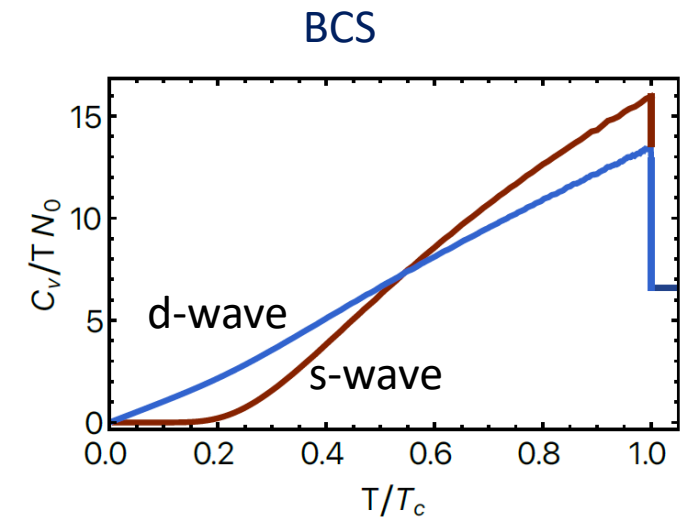


Other observables:

Specific heat: very different from BCS



- no jump at T_c , $C_v(T) \sim (T_c - T)^{1/2}$
- max near $0.8 T_c$
- “almost” finite C_v at $T \rightarrow 0$



Specific heat

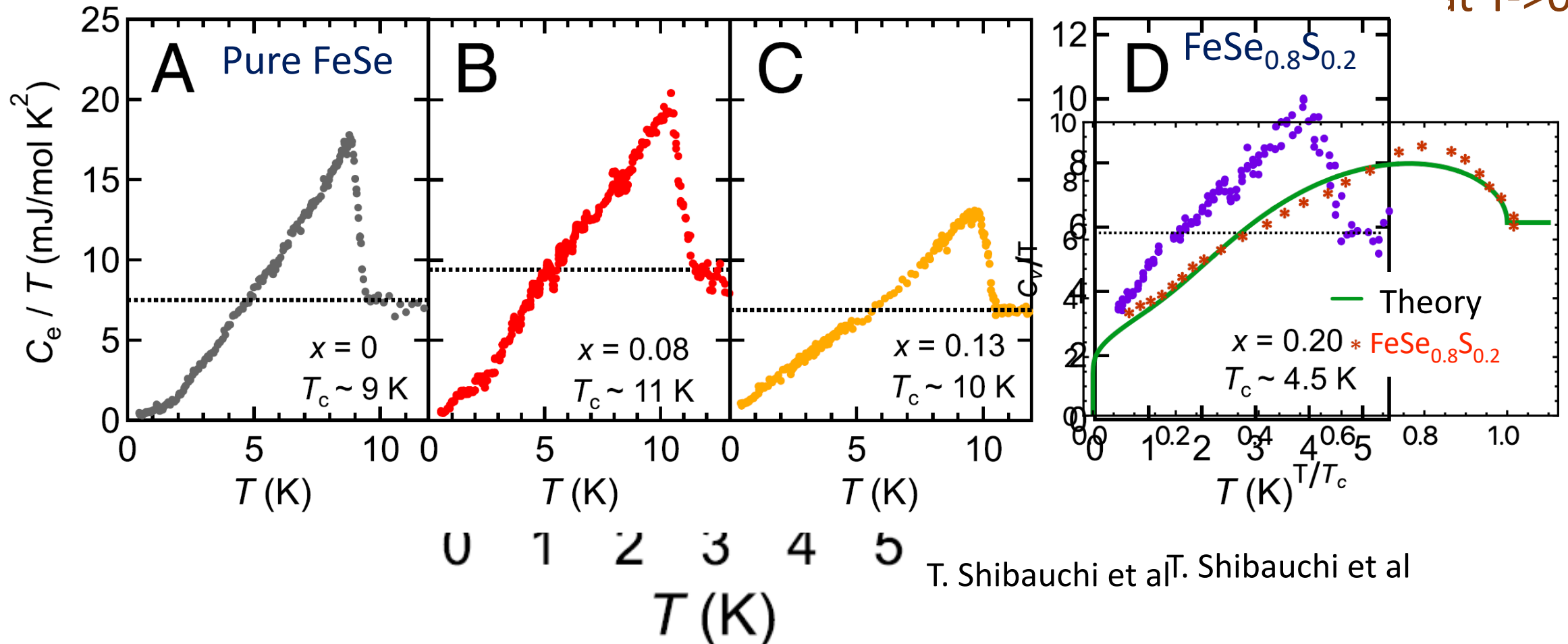


Away from a nematic QCP:

S content decreases, the system moves away from a nematic QCP

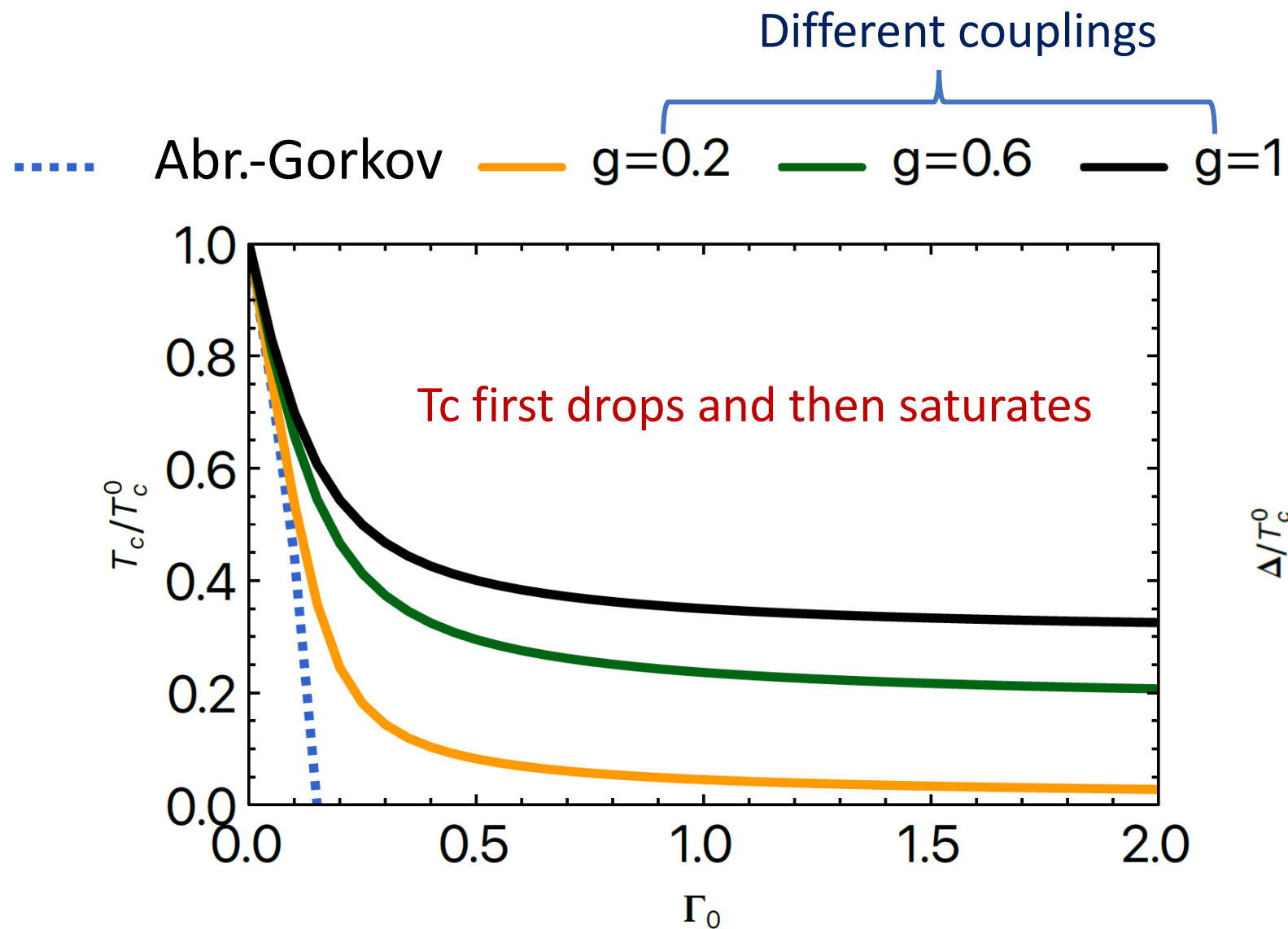
- no jump at T_c
- max near $0.8 T_c$

at $T \rightarrow 0$

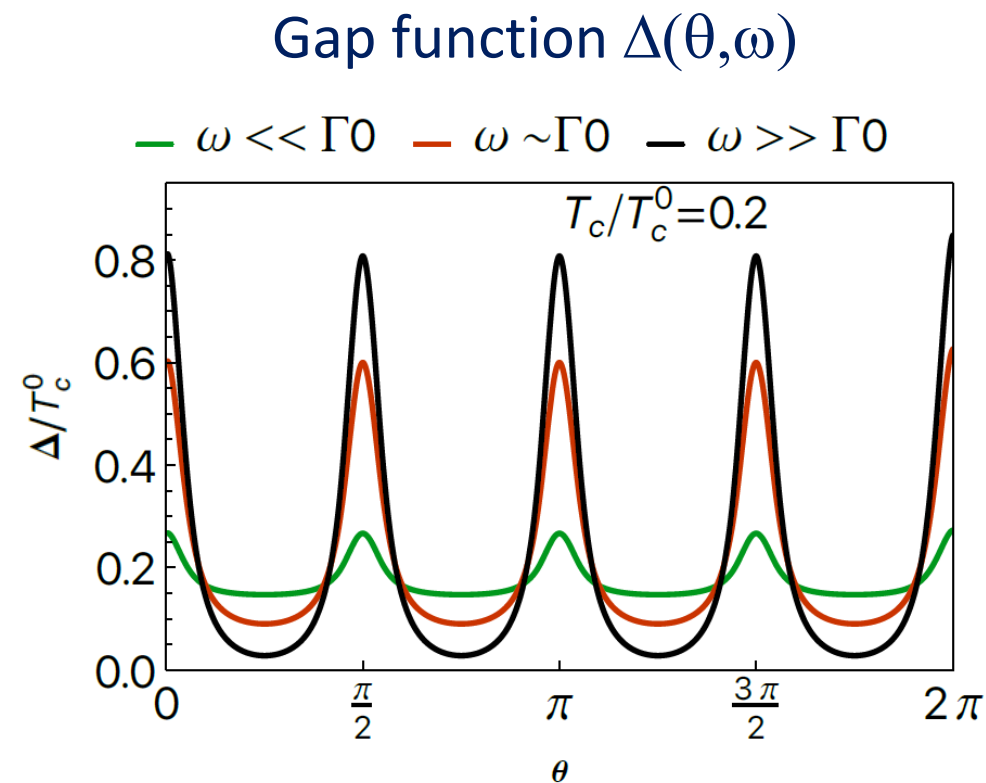


T. Shibauchi et al T. Shibauchi et al

Tc variation with impurity scattering rate

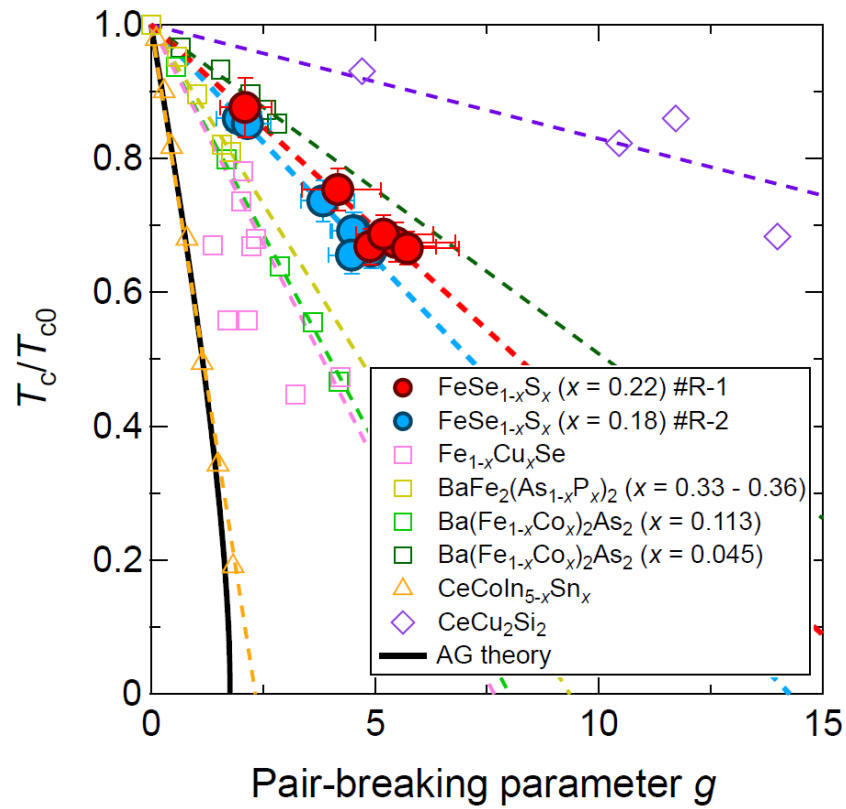


As impurity scattering increases,
gap anisotropy gets weaker



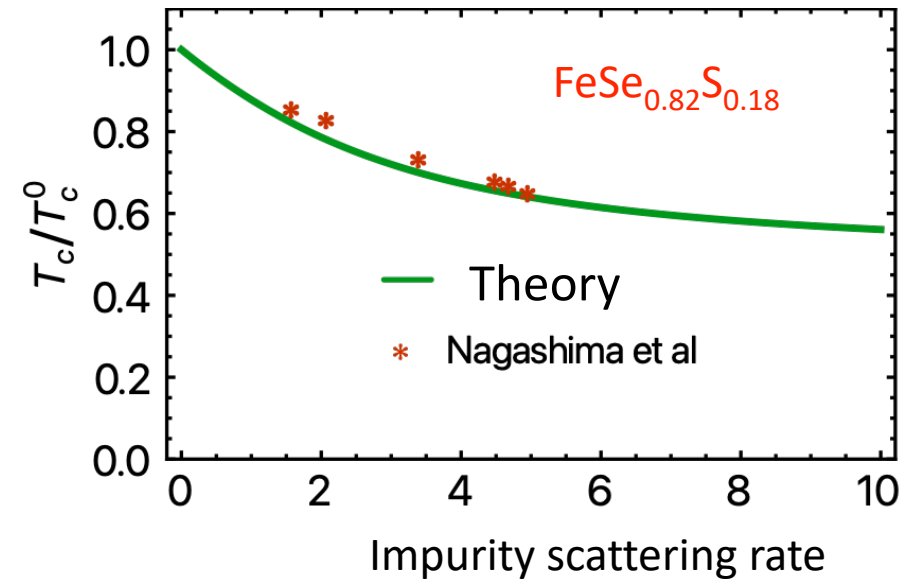
Tc variation with impurity scattering

Experiment

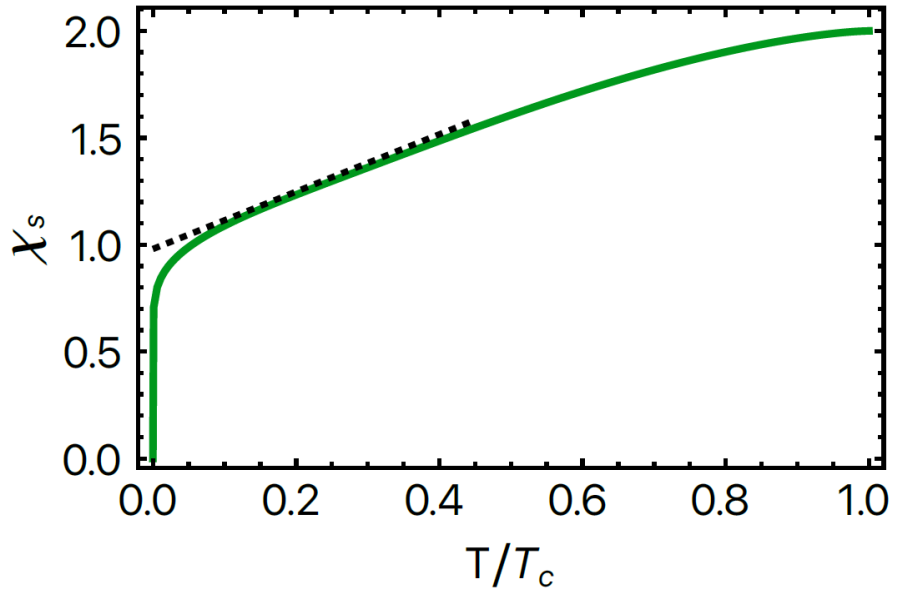


T. Nagashima et al (Shibauchi group)

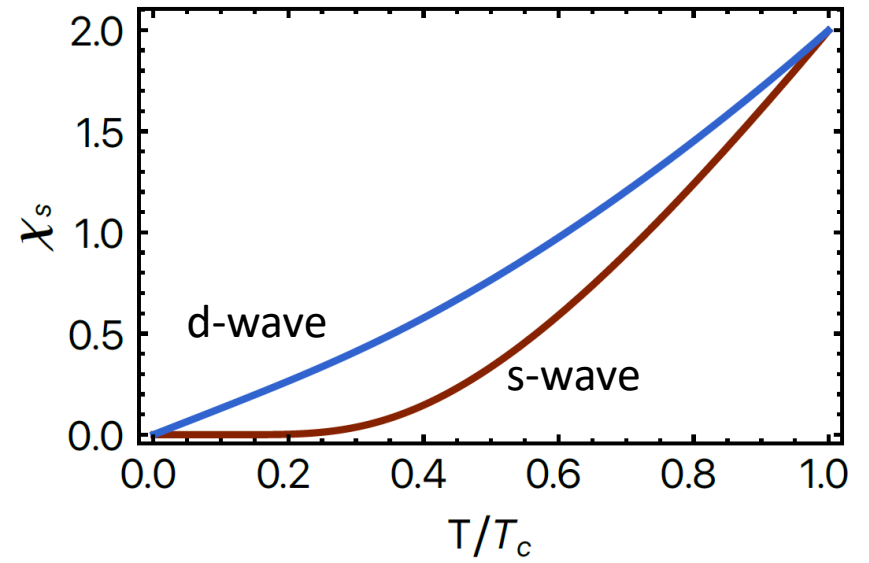
Theory



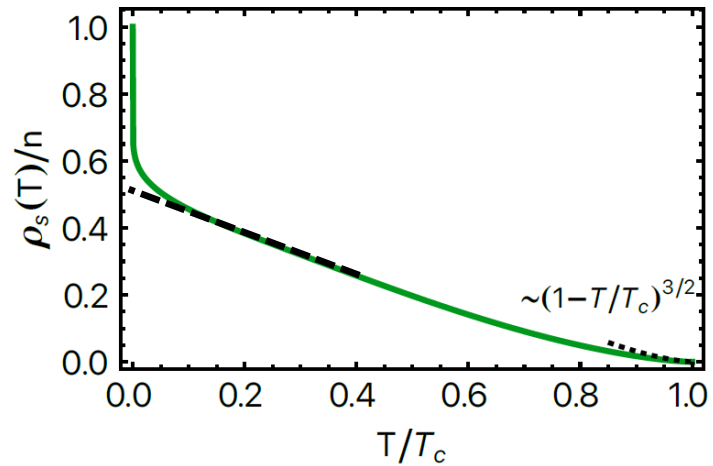
Spin susceptibility



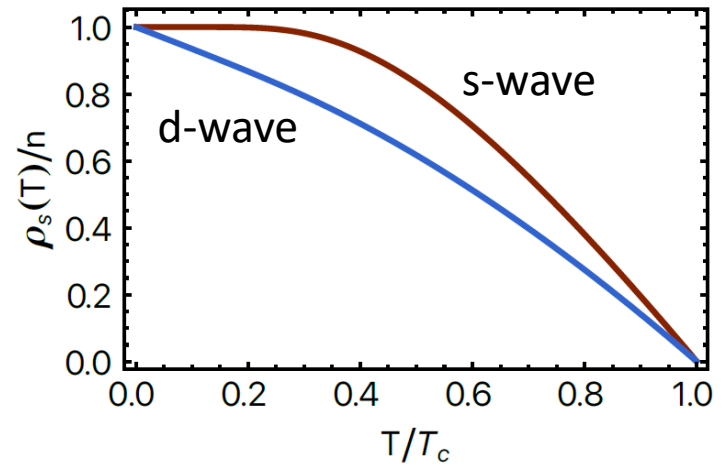
BCS



Superfluid stiffness

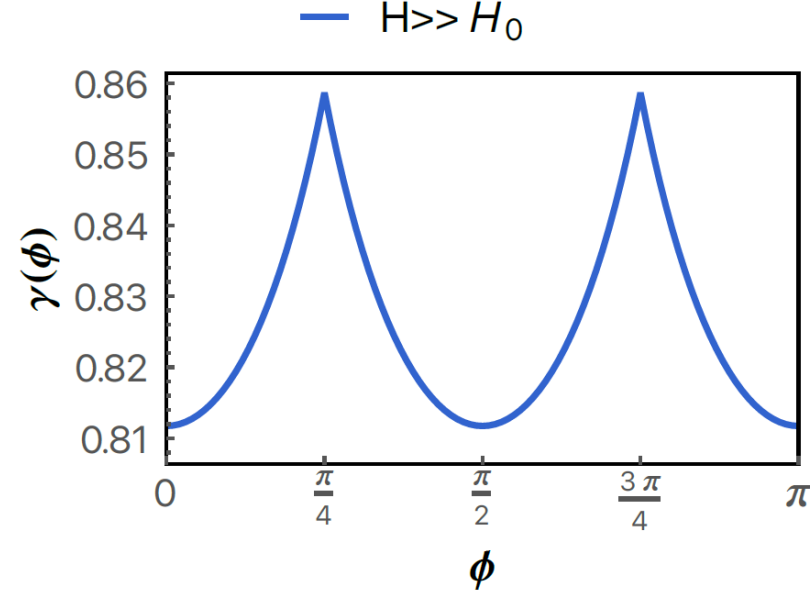
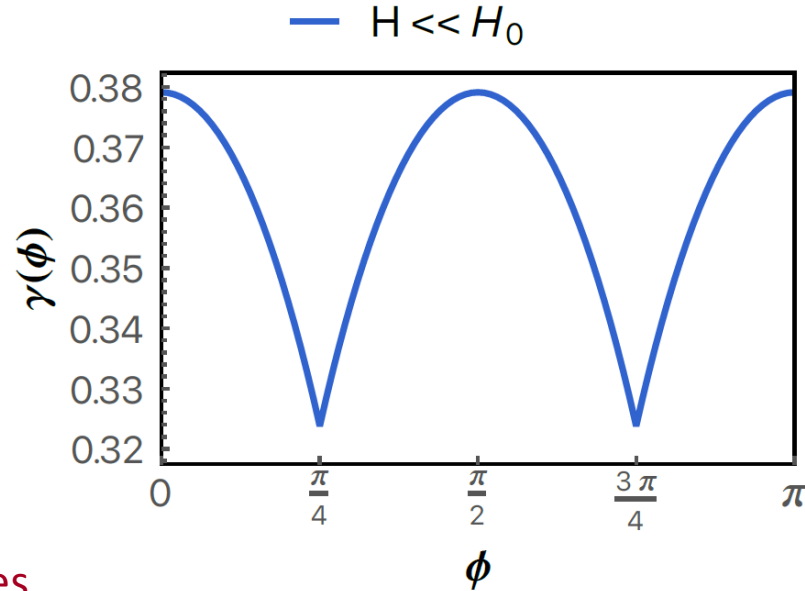


BCS



Dependence of the specific heat coefficient on the angular variation of an applied field

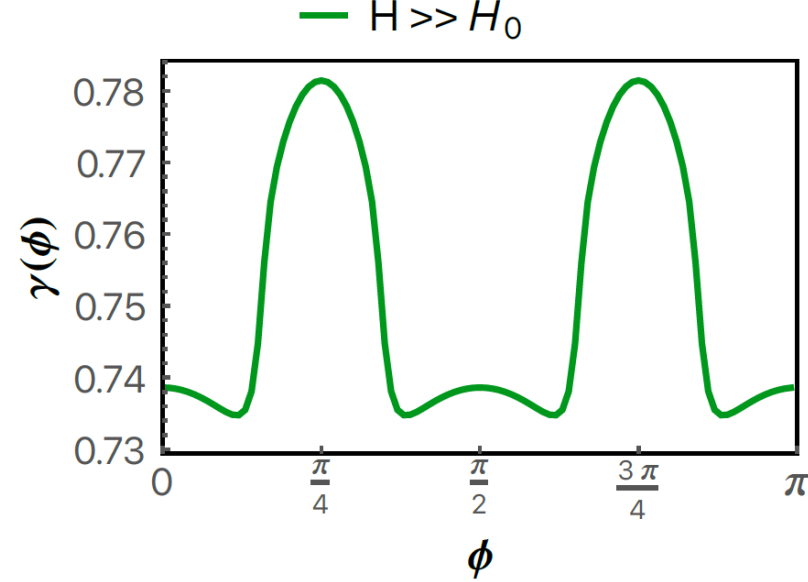
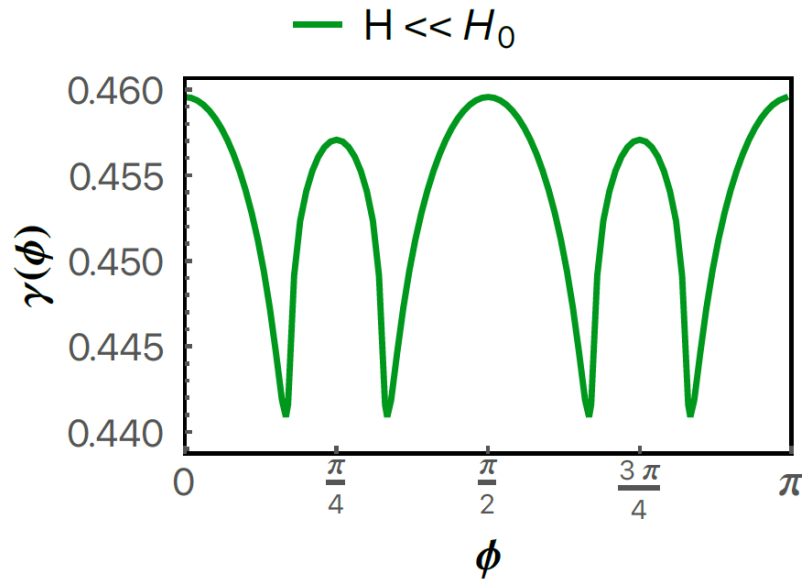
Ordinary d--wave superconductor



Vekher & Vorontsov
Eremin and A.C.

In both cases,
the gap vanishes
along the diagonal

Our case: SC mediated by nematic fluctuations

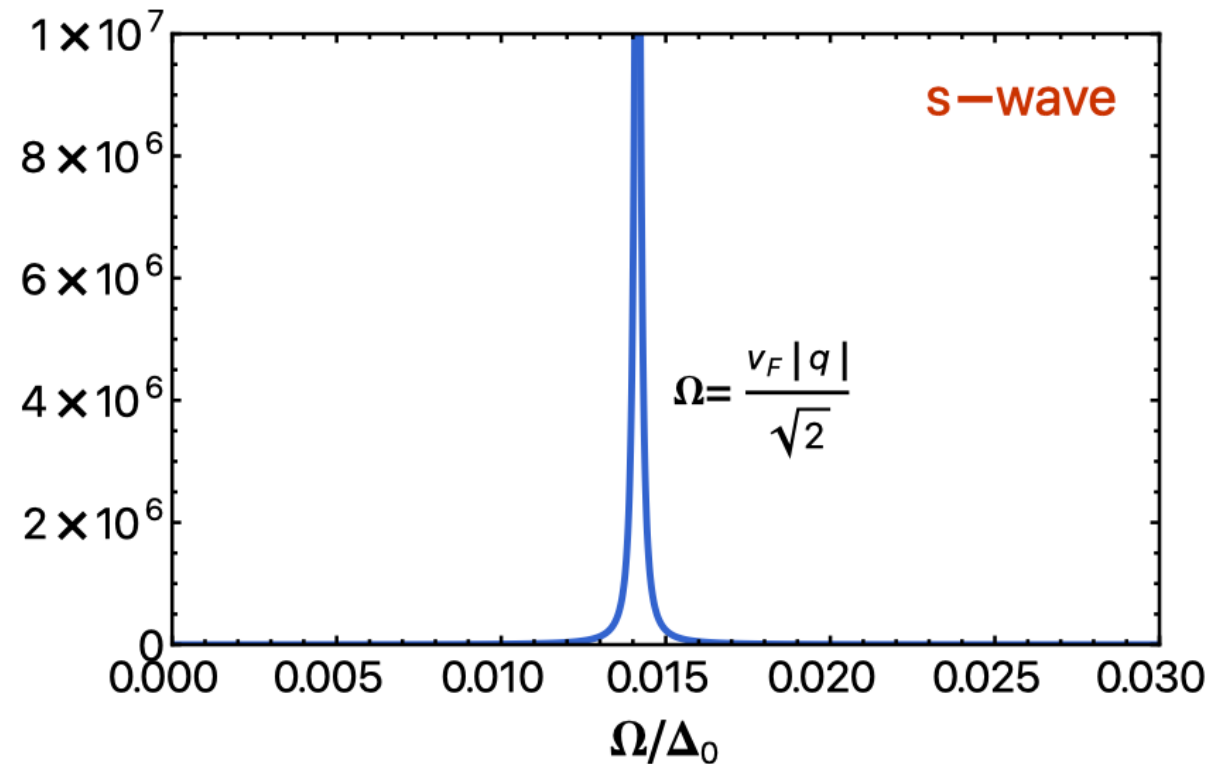


Collective modes:

Transverse pair susceptibility at various momenta
(phase fluctuations), no Coulomb.

Our case: two damped modes at
 $\Omega = v_F q \cos \theta_q$ and $\Omega = v_F q \sin \theta_q$

BCS s-wave: a single undamped
mode at $\Omega = v_F |\mathbf{q}| / \sqrt{2}$

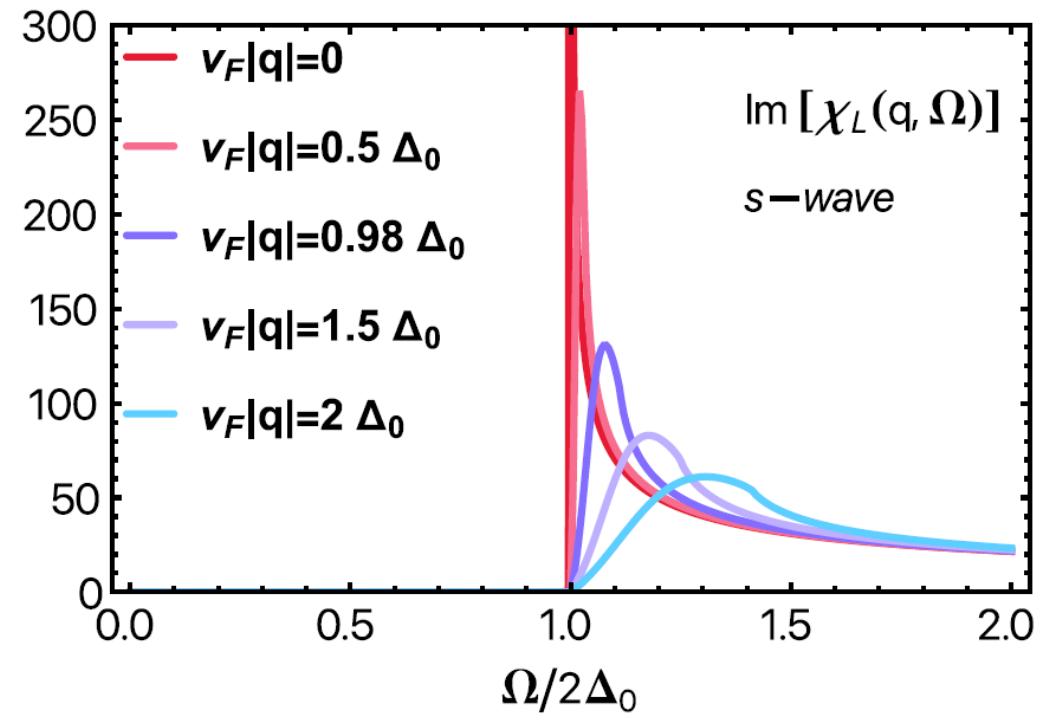


The structure of the transverse pair susceptibility
 $\chi(q, \Omega)$ depends on the direction of the momentum q

Longitudinal pair susceptibility at various momenta (longitudinal gap fluctuations)

Our case – strongly anisotropic s-wave gap

S-wave BCS superconductor



Conclusions

Superconductivity in doped FeSe near a nematic QCP
can be mediated by nematic fluctuations

- At a nematic QCP, the pairing interaction is attractive, but local in k-space: the strength of attraction depends on the position on the Fermi surface
- As the consequence, the gap opens up at T_c only at certain points, and extends along the FS as T decreases.

Features:

- Strong gap anisotropy on the hole pocket (near-nodes of the gap)
- Non-BCS behavior of the specific heat
- Impurity scattering lifts the gap in the nodal regions and reduces T_c ,
- Phase and longitudinal collective excitations are not like in a BCS SC

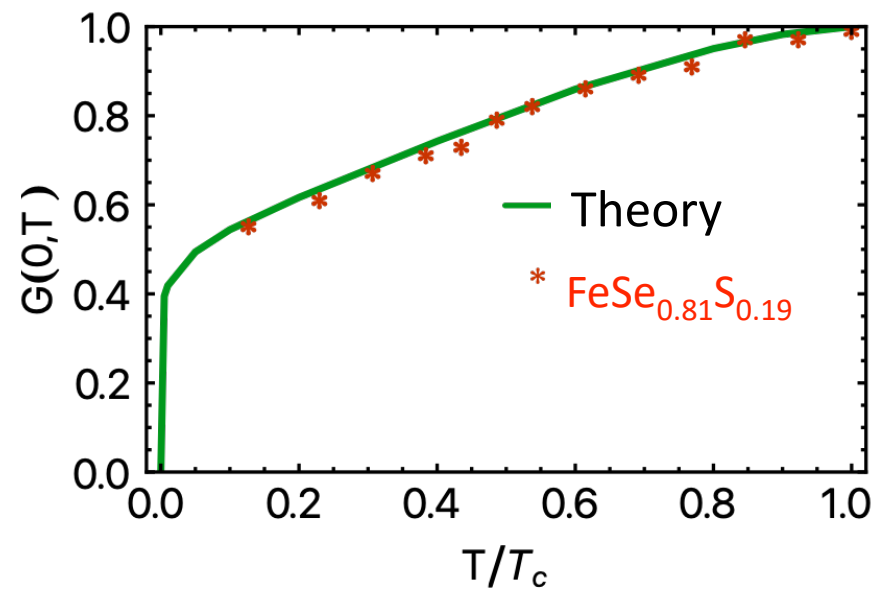
This talk



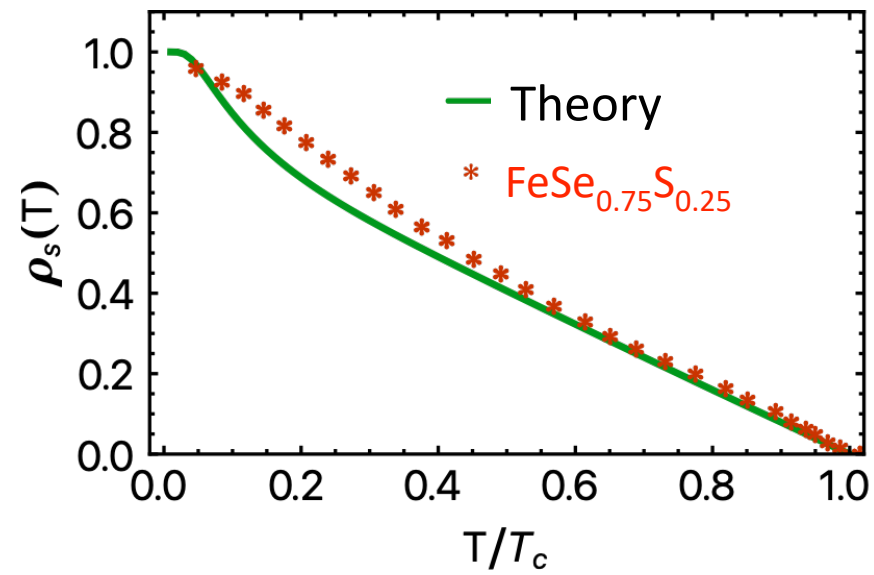
THANK YOU



Zero-bias tunneling density



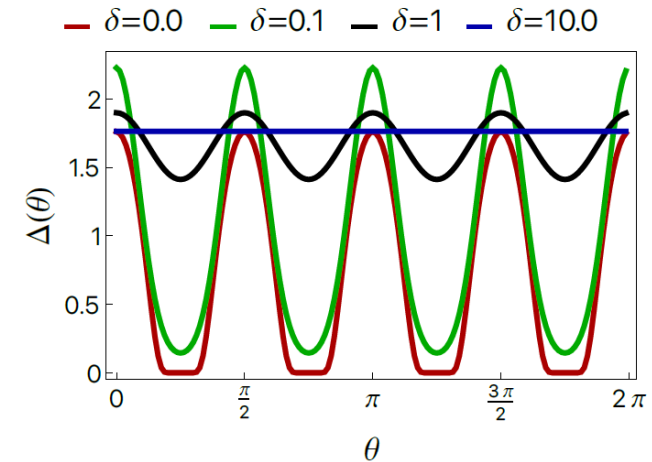
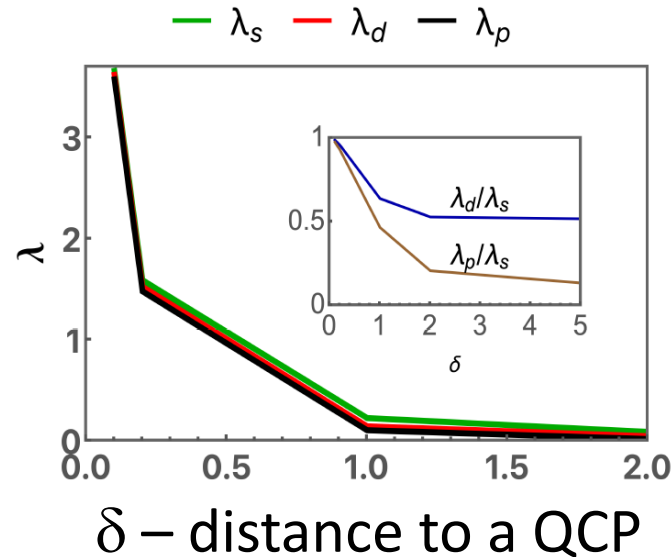
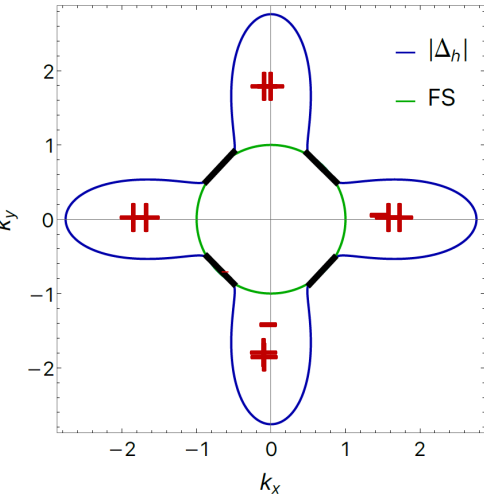
Superfluid stiffness



Experimental data:
Sato et al, Nagashima et al, Nag et al,....

Pairing symmetry

- At a nematic QCP, s-wave, d-wave, p-wave gaps are all degenerate
- Away from a QCP, or due to extra interactions, one of these states wins.
In the model of only nematic fluctuations, this is s-wave



The gap is highly anisotropic, with hot and cold regions