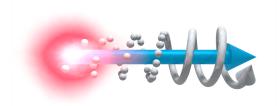
Chiral phonons for spintronics



U. Nowak University of Konstanz, Germany

Contents:

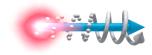
- chiral phonons and electronic angular momentum
- modelling spin lattice dynamics



Chirality and phonons







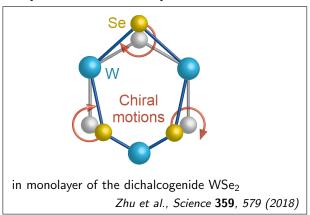


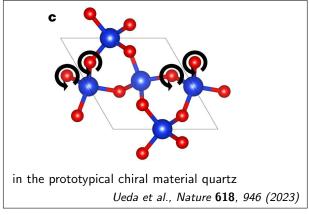
in physics: rotating and moving object

circularly polarized in a plane perpendicular to their propagation direction
 Zhang et al. PRL 115, 115502 (2015)
 Juraschek et al., Nature Physics, perspective accepted (2025)

Chiral phonons

naturally arise in non-centrosymmetric materials:





- phonon angular momentum leads to a variety of new effects:
 - significant chiral phonon magnetic moment
 - phonon Hall effect
 - phonon Faraday effect
 - phonon Zeeman effect

Juraschek et al., PRR 4, 013129 (2022)

Luo et al. Science 382, 698 (2023)

Grissonnanche et al., Nature Physics 16, 1108 (2020)

Nova et al., Nature Physics 13, 132 (2017)

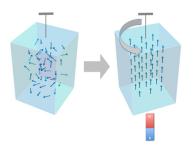
Juraschek et al., PRM 1, 014401 (2017)

Baydin et al., Phys. Rev. Lett. 128, 075901 (2022)

Einstein - De Haas effect

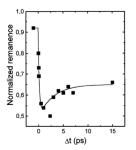
magnetic moment and angular momentum are connected

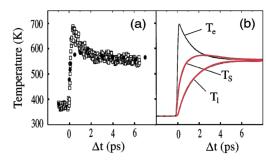
A. Einstein and W. de Haas, KNAW proceedings, **18** I, 691 (1915)
Barnett, Phys. Rev. **6**, 239 (1915)



Matsuo et al., Front. Phys. 3:54 (2015)

on fundamental time and length scales?





 magnetisation can break down on a time scale of some hundred femtoseconds

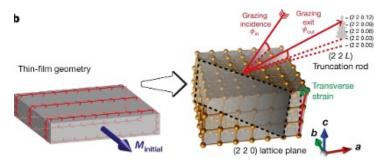
Beaurepaire et al., PRL **76**, 4250 (1996)

- phenomenological three-temperature model based on heat
- but: angular momentum conservation violated

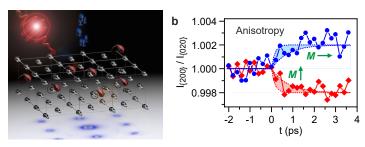
Einstein - De Haas effect

- angular momentum of phonons and the Einstein-de Haas effect
 Zhang et al., PRL 112, 085503 (2014)
- Ultrafast Einstein-de Haas effect: measured as transverse strain waves

 Ultrafast circularly polarized phonons: measured via ultrafast electron diffraction



Dornes et al., Nature 565, 209 (2019)



Tauchert et al., Nature 602, 73 (2022)

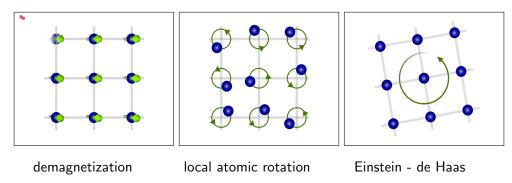


6.8.2025

Chiral phonons for spintronics

Ultrafast generation of circularly polarized phonons

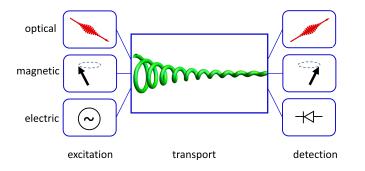
How can the lattice absorb the spin angular momentum?



New Research Unit: ChiPS



Funded by German Science Foundation: Chiral Phonons for Spintronics



Peter Baum	Visualization of chiral phonons	Uni Konstanz
	in space and time	
Ulrich Nowak	Modeling spin-lattice dynamics	Uni Konstanz
Manfred Albrecht	Advanced magnetic materials for chiral phonons	Uni Augsburg
Sangeeta Sharma	Ab-initio description of chiral- phonon and spin coupling	FU & MBI Berlin
Tobias Kampfrath	Ultrafast propagation and detec- tion of chiral phonons	FU Berlin
Sebastian Goennenwein	Chiral phonon pumping	Uni Konstanz
Silvia Viola Kusminskiy	Angular momentum transduc- tion between magnons and phon-	RWTH Aachen
	ons in patterned microstructures	
Hans Huebl	Tayloring phonon angular mo- mentum transport	WMI Garching

- fundamental length and time scales of chiral phonons
- mechanisms behind creation, detection and decay of chiral phonons
- chiral transport phenomena
- chiral phonon hybridization with other excitations
- exploit magnon-phonon interaction in tailored nanostructures
- assess potential of **phonon-based spintronics**

Chiral phonons for spintronics

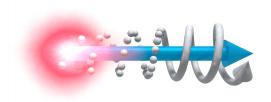
D. Schick, L. Borysenko, T. Dannegger, U. Nowak University of Konstanz, Germany

In collaboration with:

- M. Weißenhofer, Ph. Rieger, P. Oppeneer, Uppsala University
- S. Mankovsky, S. Polesya, H. Ebert, LMU München
- A. Kamra, RPTU Kaiserslautern
- L. Rosza, L. Szunyogh, Wigner Institute and BUTE Budapest

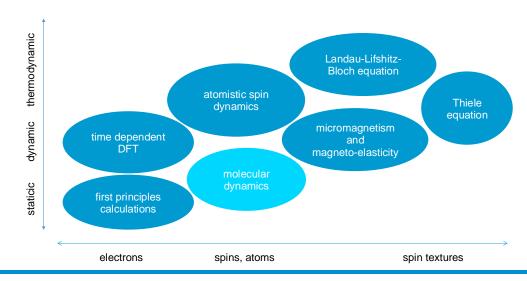
Contents:

- chiral phonons and electronic angular momentum
- modelling spin lattice dynamics



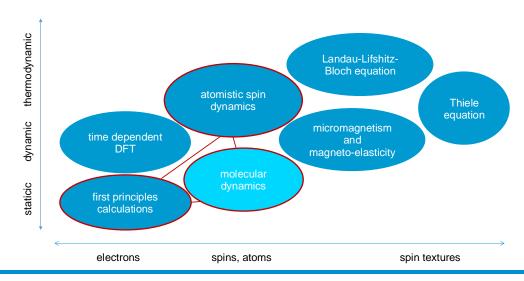
Multi-scale modelling in magnetism

- time scales from femtoseconds to years
- length scales from electronic to sample size
- temperatures from zero to above Curie temperature
- opto-magnetic effects, charge currents, laser heating



Multi-scale modelling in magnetism

- time scales from femtoseconds to years
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- temperatures from zero to above Curie temperature
- opto-magnetic effects, charge currents, laser heating



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Atomistic spin model

- Hamiltonian for spins $\underline{S}_i = \underline{\mu}_i/\mu_s$ on a given lattice:

spin model including relativistic interactions

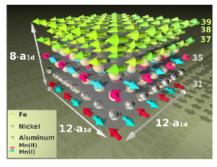
$$\mathcal{H} = -\frac{1}{2} \sum_{i,j} \mathbf{S}_i \mathbf{J}_{ij} \mathbf{S}_j - \sum_i d_i^z (S_i^z)^2 - \mathbf{B} \cdot \sum_i \mu_i \mathbf{S}_i - \sum_{i,j} \frac{\mu_0 \mu_i \mu_j}{8\pi} \frac{3(\mathbf{S}_i \cdot \mathbf{e}_{ij})(\mathbf{e}_{ij} \cdot \mathbf{S}_j) - \mathbf{S}_i \cdot \mathbf{S}_j}{r_{ij}^3}$$

exchange anisotropy external field dipole-dipole

- **tensorial exchange interactions** J_{ij} can be decomposed in:

$$\mathcal{H}_{\mathrm{ex}} = -\frac{1}{2} \sum_{i,j} J_{ij}^{\mathrm{iso}} \boldsymbol{S}_i \cdot \boldsymbol{S}_j - \frac{1}{2} \sum_{i,j} \boldsymbol{S}_i J_{ij}^{\mathrm{S}} \boldsymbol{S}_j - \frac{1}{2} \sum_{i,j} \boldsymbol{D}_{ij} \cdot (\boldsymbol{S}_i \times \boldsymbol{S}_j)$$
 isotropic exchange two-site anisotropy Dzyaloshynskii-Moriya

- different types of anisotropies . . .
- dipole-dipole interaction leads to:
 - shape anisotropy
 - domain structures
 - large numerical effort



Yanes et al., PRB 96, 064435 (2017)

Spin model parameters from first principles

exchange integrals, isotropic but orbital resolved

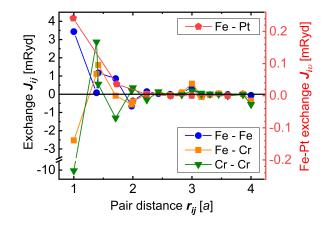
Liechtenstein et al., JMMM 67, 65 (1987)

 fully relativistic SKKR method plus spin cluster expansion for layered systems and clusters
 Szunyogh et al., PRB 83, 024401 (2011)

- calculations of **opto-magnetic effects** such as IFE

Berritta et al., PRL **117**, 137203 (2016)

John et al., Scientific Reports **7**, 4114 (2017)



Schmidt et al., PRB 102, 214436 (2020)

Gd	orbital-resolved dynamics	Frietsch et al., Nature Com. 6 8262 (2015)
Tb	orbital-resolved dynamics	Frietsch et al., Science Advances, 6 , eabb1601 (2020)
CrPt	switching with IFE in an AFM	Dannegger et al., Phys. Rev. B Lett. 104 , 060413 (2021)
Hematite	altermagnet	Dannegger et al., Phys. Rev. B 107 , 184426 (2023)

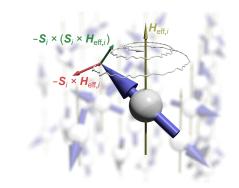
Atomistic spin dynamics

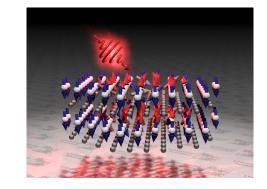
stochastic Landau-Lifshitz-Gilbert equation

$$\dot{m{S}_{\pmb{i}}} = -rac{\gamma}{(1+lpha^2)\mu_s} \, m{S}_{\pmb{i}} imes m{H}_{\pmb{i}}(t)$$
 precession $-rac{lpha\gamma}{(1+lpha^2)\mu_s} \, m{S}_{\pmb{i}} imes m{\left(m{S}_{\pmb{i}} imes m{H}_{\pmb{i}}(t)
ight)}$ dissipation

with
$$\mathbf{H}_{\mathbf{i}}(t) = -\frac{\partial \mathcal{H}}{\partial \mathbf{S}_{\mathbf{i}}} + \zeta_{\mathbf{i}}(t)$$
 and $\langle \zeta_{i\eta}(0)\zeta_{j\vartheta}(t)\rangle = \delta_{ij}\delta_{\eta\vartheta}\delta(t)2\alpha k_{\mathrm{B}}T\mu_{s}/\gamma$.

- fluctuations
- **numerical integration** for up to 10⁸ spins Lyberatos et al., J. Phys. C 5, 8911 (1993)
- statistical average in the **canonical ensemble**
- realistic dispersion relations; non-linear prozesses; critical behavior; fluctuations
- © electrons and phonons only as heat-bath with coupling constant α ; classical approximation; large numerical effort





Spin-lattice dynamics

Hamiltonian for spin and lattice degrees of freedom

$$\mathcal{H}(\boldsymbol{S}, \boldsymbol{r}, \boldsymbol{\rho}) = \underbrace{\sum_{i,j} J(\boldsymbol{S}_i, \boldsymbol{S}_j)}_{\text{spin-lattice}} + \underbrace{\sum_{i} \frac{\boldsymbol{\rho}_i^2}{2m_i} + \sum_{i,j} V(\boldsymbol{r}_{ij})}_{\text{lattice}}$$

equations of motion

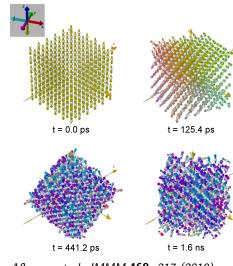
$$\dot{\boldsymbol{r}}_i = \frac{\partial \mathcal{H}}{\partial \boldsymbol{p}_i}, \quad \dot{\boldsymbol{p}}_i = -\frac{\partial \mathcal{H}}{\partial \boldsymbol{r}_i}, \quad \dot{\boldsymbol{S}}_i = \frac{\gamma}{\mu_s} \boldsymbol{S}_i \times \frac{\partial \mathcal{H}}{\partial \boldsymbol{S}_i}$$

Strungaru et al., PRB **103**, 024429 (2021) Hellsvik et al., PRB **99**, 104302 (2019)

- **no heat bath** → microcanonic ensemble
- open problems:

6.8.2025

- calculation of spin-lattice coupling parameters from first-principles
- correct expressions for spin-lattice coupling
- conservation of energy, momentum, angular momentum



Aßmann et al, JMMM 469, 217 (2019)

Chiral phonons for spintronics

Spin-lattice coupling from first principles

expansion of relativistic Heisenberg Hamiltonian in terms of displacements $u_i = r_i - r_i^{eq}$

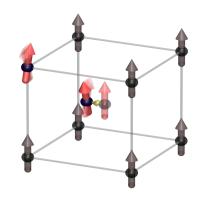
spin-lattice Hamiltonian

$$\mathcal{H}^{\mathrm{SLC}} = \sum_{ij,\alpha\beta} J_{ij}^{\alpha\beta}(\mathbf{r}_{k}) S_{i}^{\alpha} S_{j}^{\beta} = \sum_{ij,\alpha\beta} J_{ij}^{\alpha\beta}(\mathbf{r}^{\mathrm{eq}}) S_{i}^{\alpha} S_{j}^{\beta} + \sum_{ijk,\alpha\beta\mu} J_{ijk}^{\alpha\beta\mu}(\mathbf{r}^{\mathrm{eq}}) S_{i}^{\alpha} S_{j}^{\beta} u_{k}^{\mu} + \dots$$

$$\underbrace{}_{\mathrm{spin-spin}}^{\mathrm{spin-spin}} \underbrace{}_{\mathrm{interactions}}^{\mathrm{spin-spin-lattice}} \underbrace{}_{\mathrm$$

spin-lattice coupling tensors $J_{ijk}^{lphaeta\mu}$ can be calculated via:

- super cell method Hellsvik et al., PRB **99**, 104302 (2019)
- extension of Liechtenstein formula
 Mankovsky et al., PRL 129, 067202 (2022)
- embedded cluster method
 Lange et al., PRB 107, 115176 (2023)



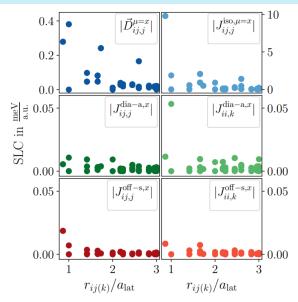
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Spin-lattice coupling parameters for bcc iron from first principles

spin-lattice Hamiltonian

$$\mathcal{H}_{sl} = \sum_{i,j,\alpha,\beta} J_{ij}^{\alpha\beta} S_i^{\alpha} S_j^{\beta} + \sum_{i,j,\alpha,\beta} \sum_{k,\mu} J_{ij,k}^{\alpha\beta,\mu} S_i^{\alpha} S_j^{\beta} u_k^{\mu} + \sum_{i,j,\alpha,\beta} \sum_{k,l,\mu,\nu} J_{ij,kl}^{\alpha\beta,\mu\nu} S_i^{\alpha} S_j^{\beta} u_k^{\mu} u_l^{\nu}$$

- example: bcc Fe
- fully-relativistic scheme treats changes of spin configuration and atomic positions equally
- even in inversion symmetric lattice leading term for angular momentum transfer is DMI-type interaction
- beyond magneto-elastic theory



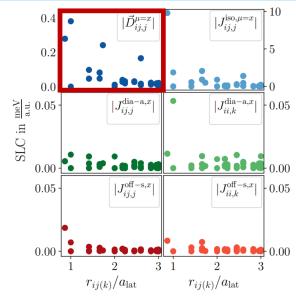
Mankovsky et al., PRL 129, 067202 (2022)

Spin-lattice coupling parameters for bcc iron from first principles

spin-lattice Hamiltonian

$$\mathcal{H}_{sl} = \sum_{i,j,\alpha,\beta} J_{ij}^{\alpha\beta} S_i^{\alpha} S_j^{\beta} + \sum_{i,j,\alpha,\beta} \sum_{k,\mu} J_{ij,k}^{\alpha\beta,\mu} S_i^{\alpha} S_j^{\beta} u_k^{\mu} + \sum_{i,j,\alpha,\beta} \sum_{k,l,\mu,\nu} J_{ij,kl}^{\alpha\beta,\mu\nu} S_i^{\alpha} S_j^{\beta} u_k^{\mu} u_l^{\nu}$$

- example: bcc Fe
- fully-relativistic scheme treats changes of spin configuration and atomic positions equally
- even in inversion symmetric lattice leading term for angular momentum transfer is DMI-type interaction
- beyond magneto-elastic theory



Mankovsky et al., PRL 129, 067202 (2022)

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Towards combined spin-lattice dynamics

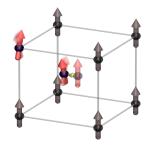
Full spin-lattice Hamiltonian

$$\mathcal{H} = \sum_{ij,\alpha\beta} J_{ij}^{\alpha\beta} S_i^{\alpha} S_j^{\beta} + \sum_{ijk,\alpha\beta\mu} J_{ijk}^{\alpha\beta\mu} S_i^{\alpha} S_j^{\beta} u_k^{\mu} + \sum_i \frac{\boldsymbol{p}_i^2}{2m_i} + \sum_{ij} V(\boldsymbol{r}_{ij})$$

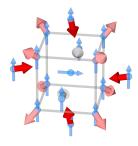
resulting equations of motion:

$$\dot{\boldsymbol{S}}_{i} = rac{\gamma}{\mu_{\mathrm{s}}} \boldsymbol{S}_{i} imes rac{\partial \mathcal{H}}{\partial \boldsymbol{S}_{i}} = -rac{\gamma}{\mu_{\mathrm{s}}} \boldsymbol{S}_{i} imes \left(\boldsymbol{H}_{i}^{\mathrm{SS}} + \boldsymbol{H}_{i}^{\mathrm{SLC}} \right)$$
 $\dot{\boldsymbol{r}}_{i} = rac{oldsymbol{p}_{i}}{m_{i}}$
 $\dot{oldsymbol{p}}_{i} = -rac{\partial \mathcal{H}}{\partial oldsymbol{r}_{i}} = oldsymbol{F}_{i}^{\mathrm{lattice}} + oldsymbol{F}_{i}^{\mathrm{SLC}}$

- spin lattice coupling breaks rotational symmetry
- ⇒ total angular momentum not conserved! (see also *Melcher*, *PRL* **25**, *1201* (*1970*) for corresponding magneto-elastic theory)
 - what about the chiral phonon magnetic moment? $S_i \cdot (u_i \times p_i)$



SLC-fields



SLC-forces

6.8.2025

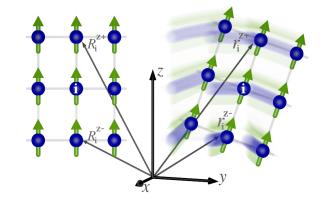
Angular momentum conserving formulation of spin-lattice dynamics

replacing components of spins and displacements with **projection onto local lattice orientation**

$$oldsymbol{e}_{i}^{lpha}=rac{oldsymbol{r}_{i}^{lpha+}-oldsymbol{r}_{i}^{lpha-}}{|oldsymbol{r}_{i}^{lpha+}-oldsymbol{r}_{i}^{lpha-}|}$$

Resulting spin-lattice Hamiltonian

$$egin{aligned} \mathcal{H} &= \sum_{ij,lphaeta} J_{ij}^{lphaeta}(oldsymbol{S}_i\cdotoldsymbol{e}_i^lpha)(oldsymbol{S}_j\cdotoldsymbol{e}_j^eta) \ &+ \sum_{ijk,lphaeta\mu} J_{ijk}^{lphaeta\mu}(oldsymbol{S}_i\cdotoldsymbol{e}_i^lpha)(oldsymbol{S}_j\cdotoldsymbol{e}_j^eta)(oldsymbol{u}_k\cdotoldsymbol{e}_k^\mu) \ &+ \sum_{i} rac{oldsymbol{p}_i^2}{2m_i} + \sum_{ii} V(oldsymbol{r}_{ij}) \end{aligned}$$



e.g. direction of easy axis is defined by the lattice not via the lab frame

$$d_z \sum_i (S_i^z)^2 \rightarrow d_z \sum_i (\boldsymbol{S}_i \cdot \boldsymbol{e}_i^z)^2$$

- this Hamiltonian is based on difference vectors and scalar products
- ⇒ it is rotational and translational invariant
- ⇒ momentum and angular momentum are conserved in the total system!

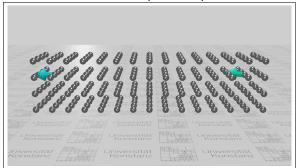
Weißenhofer et al., PRB 108, L060404 (2023)

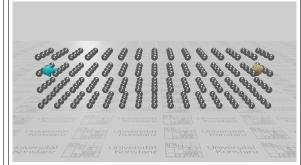


Spin-spin interaction via phonons



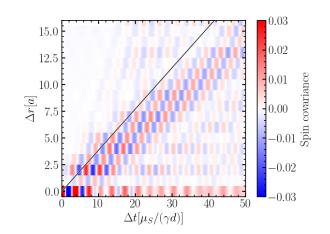






Two parallel and magnetically uncoupled spins in a phonon bath couple via exchange of angular momentum

Spin-spin interaction via phonons

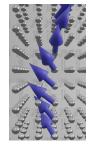


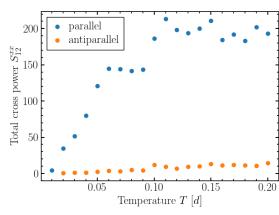
correlation quantified via cross covariance of spin components:

$$K^{\mu
u}_{ij}(\Delta t) = \langle S^{\mu}_i(t+\Delta t)S^{
u}_j(t)
angle - \langle S^{\mu}_i
angle\langle S^{
u}_j
angle$$

and cross power spectral density (CPSD):

$$P_{ij}^{\mu\nu} = \int_0^\infty \left| \tilde{S}_{ij}^{\mu\nu}(f) \right|^2 \, \mathrm{d}f$$

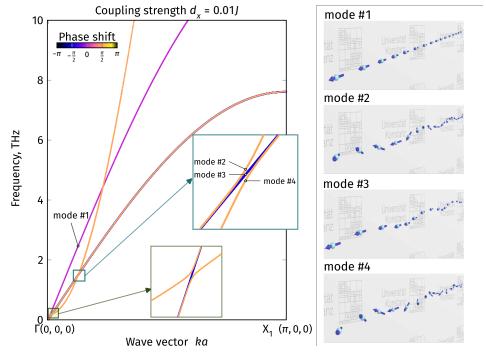




Schick et al., in preparation see also: Yokoyama, JPSJ **93**, 123705 (2024)

6.8.2025 Chiral

Chirality-selective magnon-phonon dispersion: ferromagnets

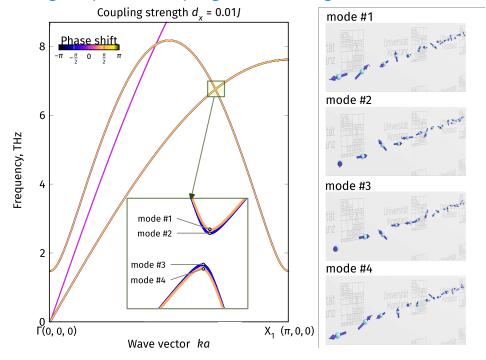


- avoided crossings, magnon polarons (Li et al., APL Materials 9, 060902 (2021)) Borysenko et al., in preparation
- degeneracy lifted: transverse phonons turn into chiral, only one of the two chiral phonon modes affected

6.8.2025

Chiral phonons for spintronics

Chirality-selective magnon-phonon coupling: antiferromagnets

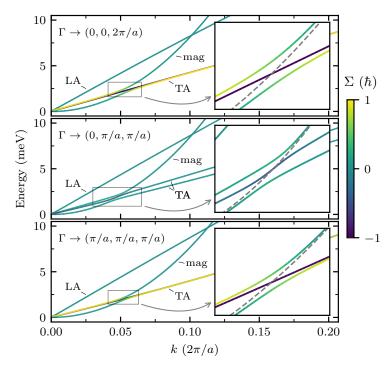


- avoided crossings, hybridized quasiparticles
- degeneracy lifted: both chiral phonon modes affected

Borysenko et al., in preparation

Chiral phonons from chirality-selective magnon-phonon coupling

- coupled magnon-phonon bands in bcc Fe from first principles via Holstein-Primakoff transformation
- grey dashed lines are the bare magnon energies
- colors encode the phonon chirality
- transverse modes are degenerate for $\Gamma o (0,0,2\pi/a)$ and $\Gamma o (\pi/a,\pi/a,\pi/a)$
- degeneracy lifted: only one of the two chiral phonon modes affected



Weißenhofer et al., arXiv:2411.03879 (2024)



Thanks to

My group in Konstanz:



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Universität Konstanz





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- M. Albrecht, Universität Augsburg
- H. Ebert, LMU München
- T. Kampfrath, FU Berlin
- A. Kamra, University of Madrid
- T. Kurihara, Tokyo University
- M. Kläui, Universität Mainz
- R. Mondal, IIOT, Dhanbad
- P. Oppeneer, M. Weißenhofer, Uppsala University
- L. Rózsa, L. Szunyogh, L. Udvardi, Budapest University

SCCKN

Scientific Compute Cluster Konstanz



Research Unit ChiPS



Fluctuations



Schwerpunktprogramm Skyrmionics

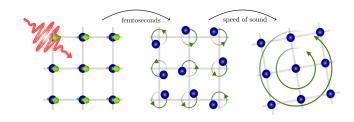
Summary:

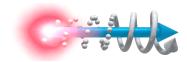
- ultrafast transfer of spin angular momentum into the lattice

- polarized phonons absorb the spin angular momentum in Ni
- chiral phonons for spintronics

- modelling spin-lattice dynamics

- relativistic spin-lattice coupling parameters from first principles
- rotationally invariant formulation of spin-lattice Hamiltonian
- phonon mediated spin spin interaction
- chirality-selective magnon-phonon coupling







6.8.2025

Chiral phonons for spintronics Ulrich Nowak, Universität Konstanz 26