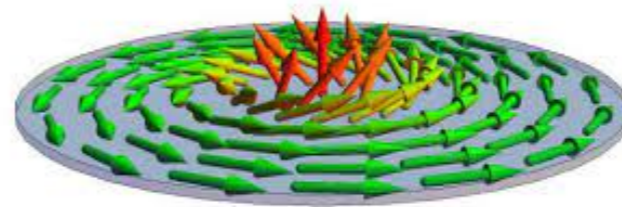
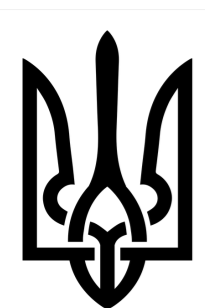


# Energy, geometry, and topology of collective magnetic dynamics

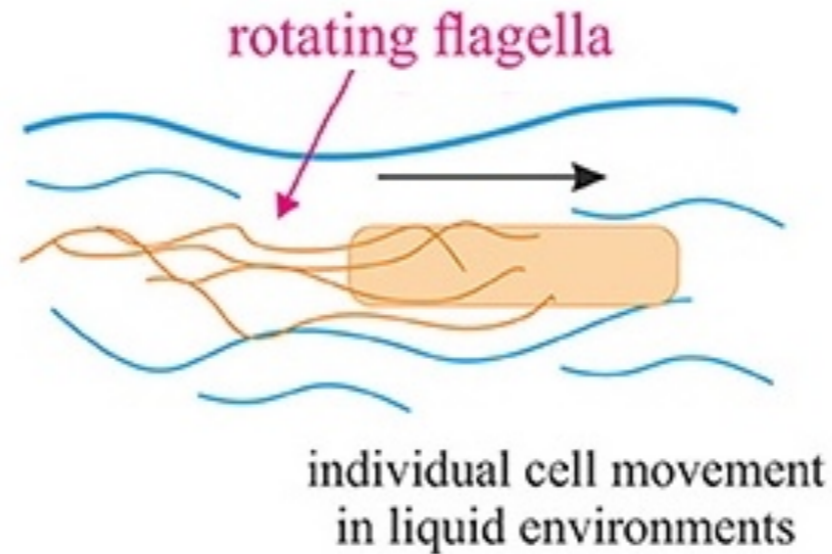


Yaroslav Tserkovnyak ([UCLA](#))

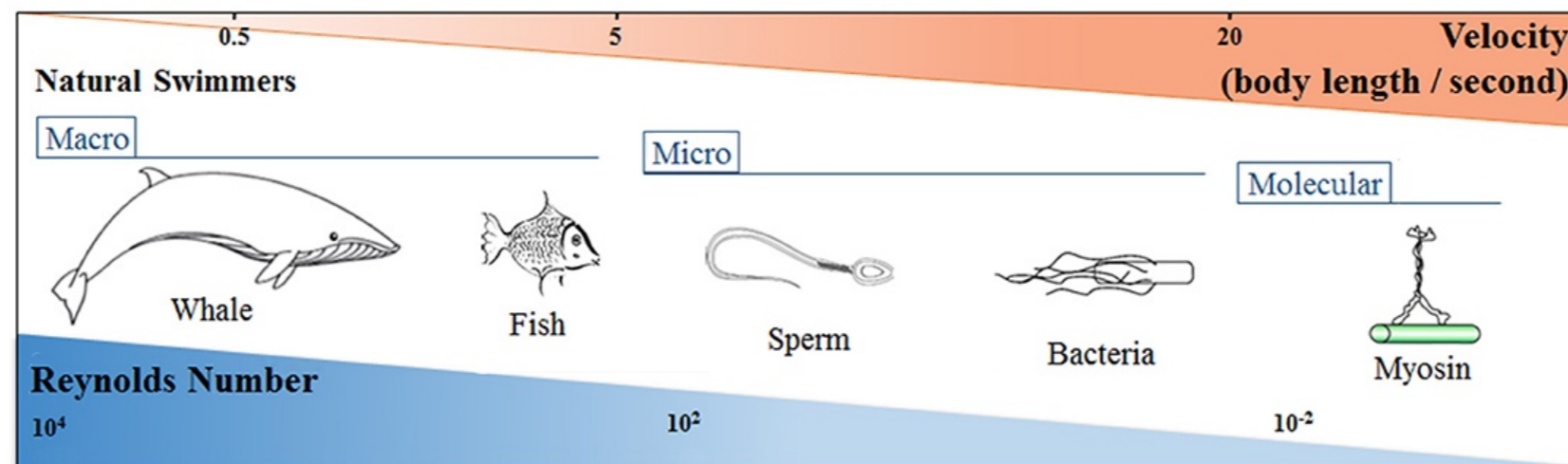
So Takei ([CUNY](#)), Se Kwon Kim ([KAIST](#)), Hector Ochoa ([Donostia](#)),  
Ricardo Zarzuela ([Mainz](#)), Pramey Upadhyaya ([Purdue](#)), Benedetta Flebus ([Boston](#)), Ji Zou ([Basel](#)),  
Suzy Zhang ([Max Planck/Dresden](#)), Chau Dao ([UCLA](#)), and Eric Kleinherbers ([UCLA](#))



# Bacterial motility

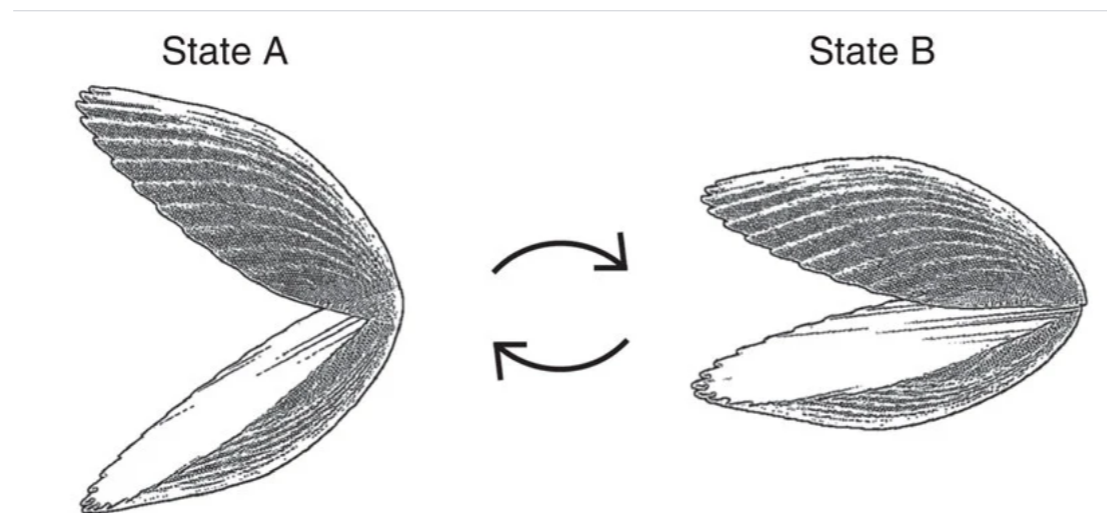


- Bacterial ability to swim using metabolic energy
- Each single flagellum (helical appendage) has a rotary motor at the base that can turn clock- or anticlockwise
- Small Reynolds number: Viscosity-dominated hydrodynamics; thus, continuous transduction of chemical into mechanical energy

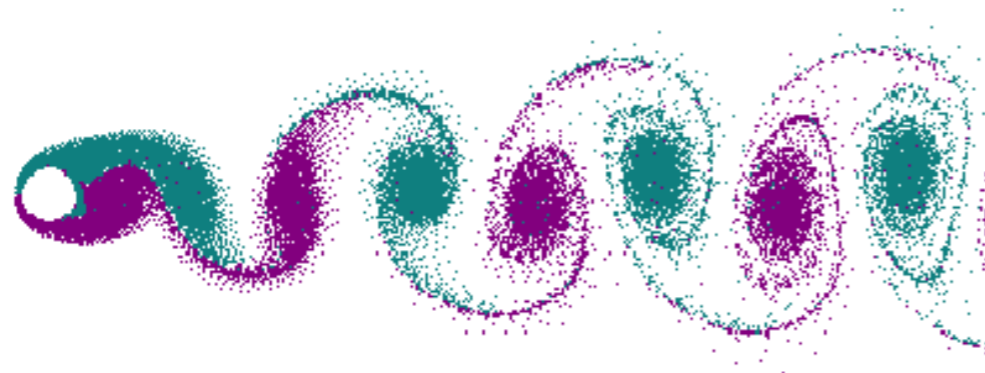


# Macroscopic mobility strategies

- Inertia dominated: Typically imparting momentum to the fluid by discrete events, such as vortex shedding, with inertial coasting in between

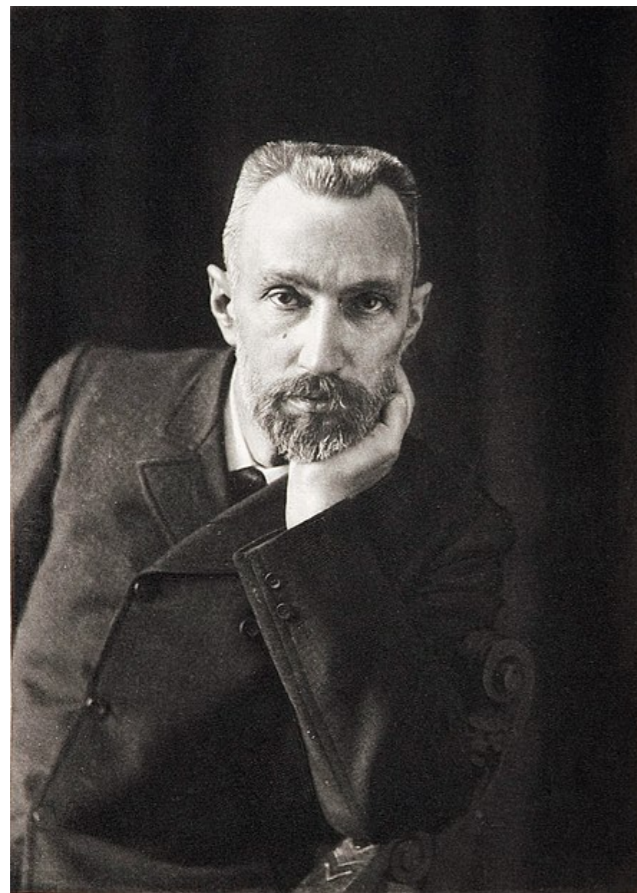


E. M. Purcell: "Fast or slow, it exactly retraces its trajectory and it's back where it started." [without inertia]



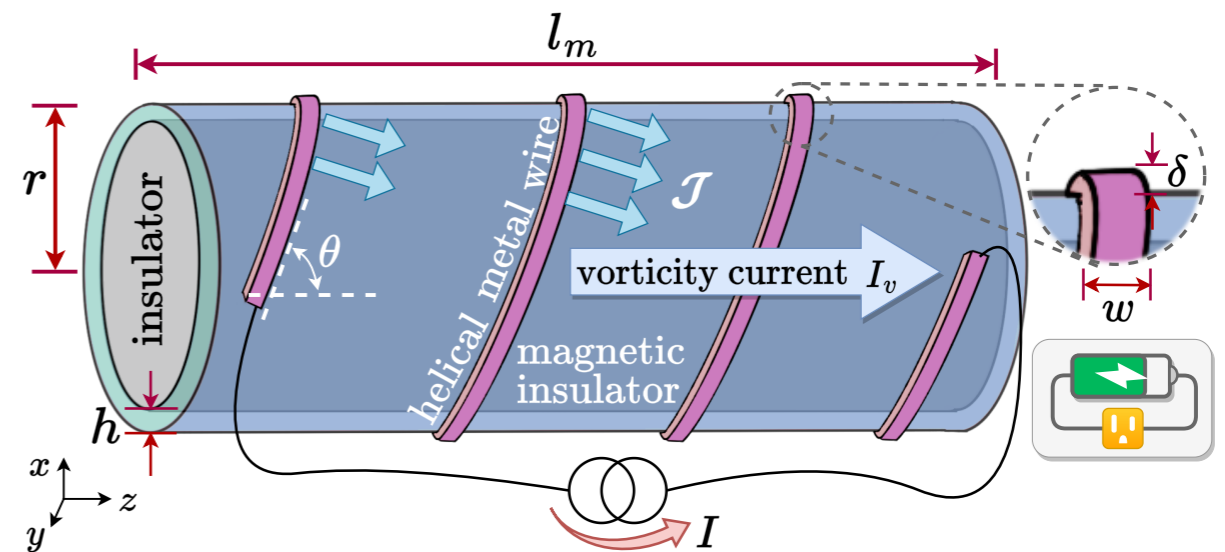
# From evolution to “intelligent” design

- Conversion of featureless electrochemical free energy into dynamic macroscopic configurations
- Geometry is exploited to enable desired process via the reduction of structural symmetries



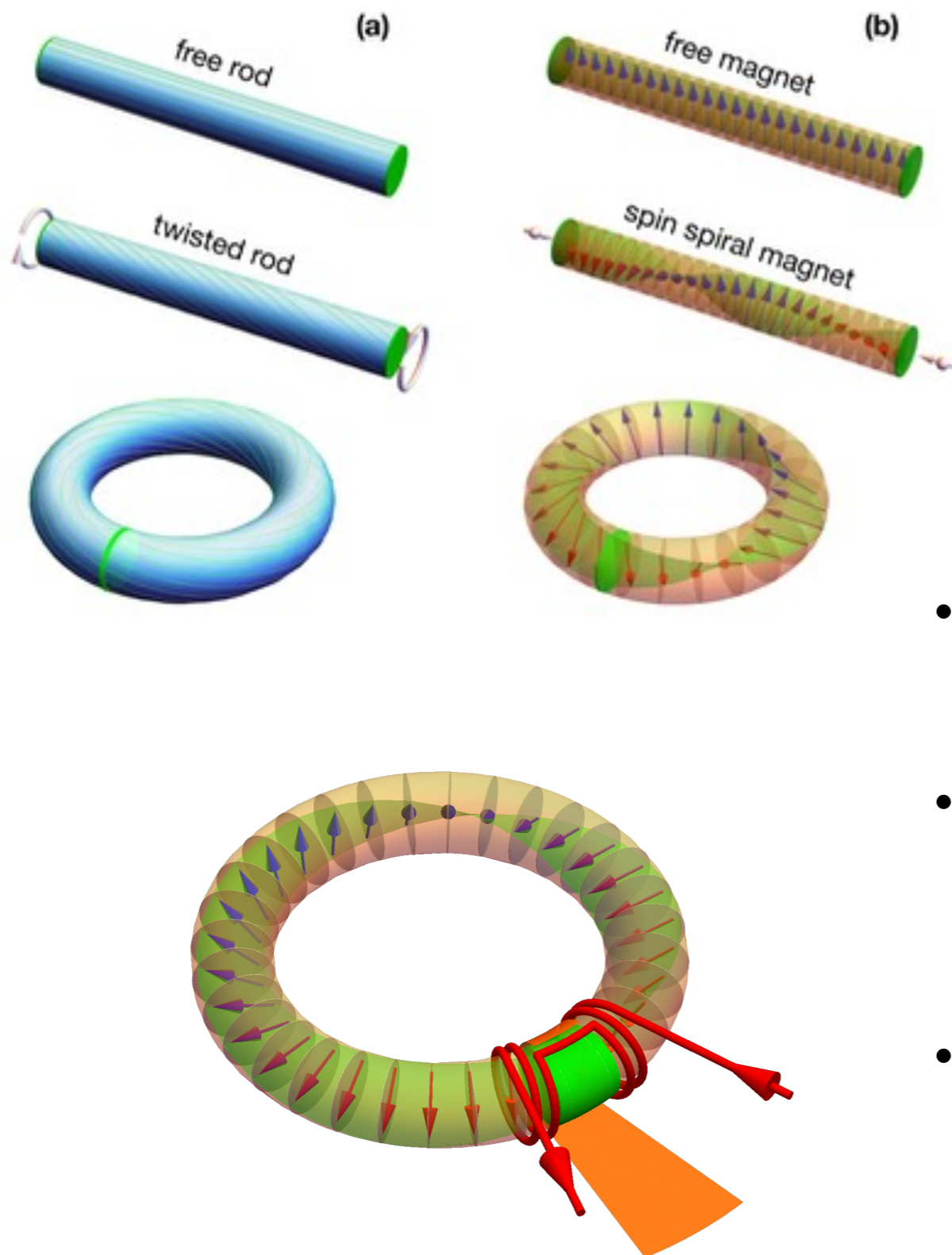
Pierre Curie:

“Asymmetry creates the phenomenon”



- Enable magnetic transport phenomena via geometry
- Utilize topology of collective spin textures to transduce and store free energy
- Interface with (thermo)electric input/output

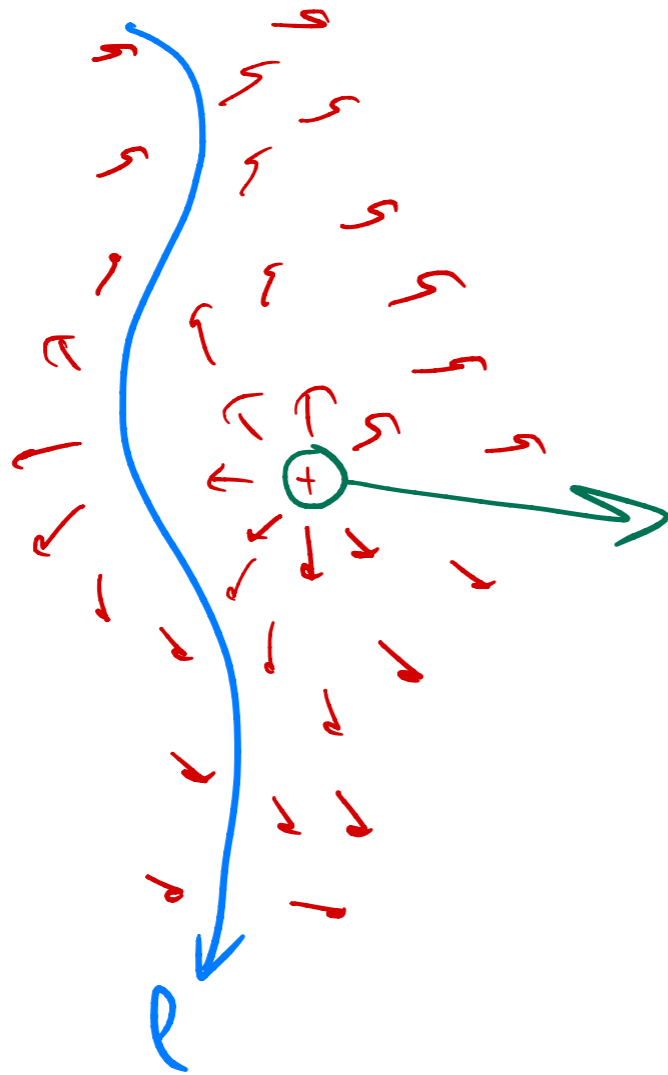
# The energy-storage concept



- How to inject/extract such winding in practice?
- Controlling symmetries: Heterostructure design vs curved geometry
- Thermodynamic efficiency of the cross-talk between winding dynamics and electricity?

# An appealing approach

---



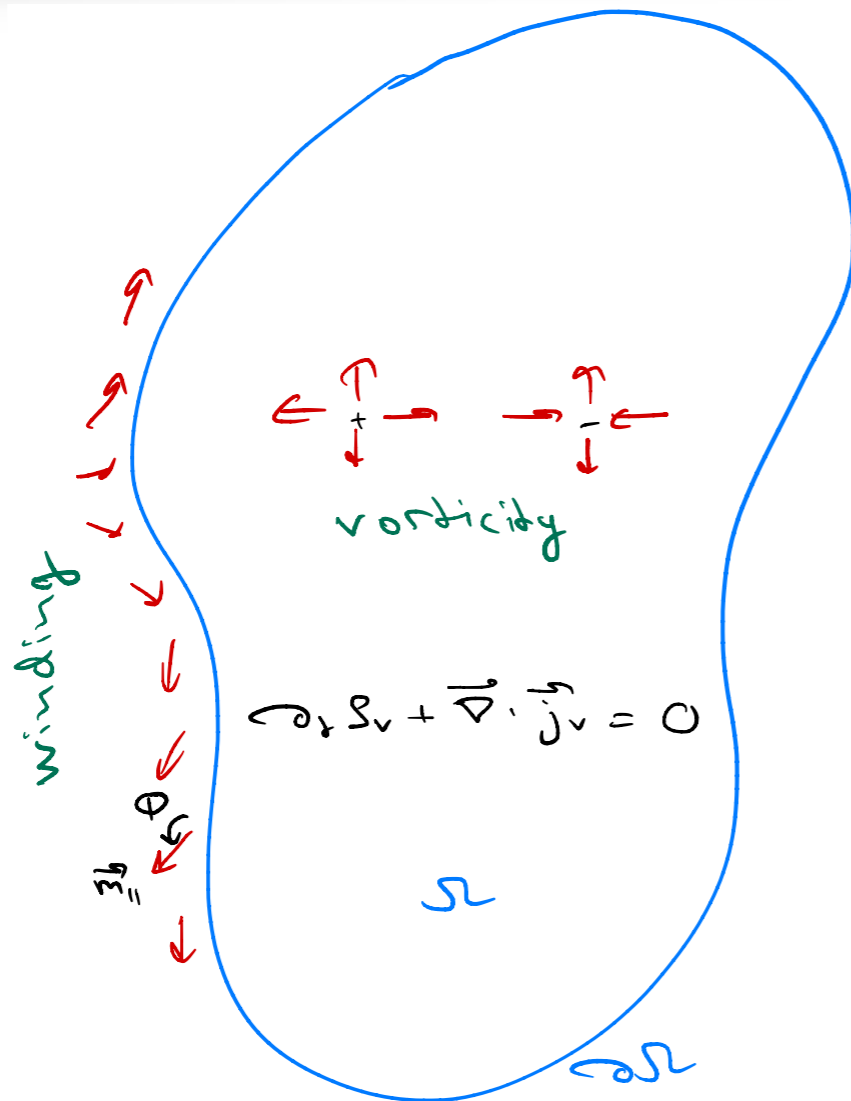
vorticity flow  
transverse to  $\rho$

$\Rightarrow$

winding change  
along  $\rho$

If we establish practical means to push vorticity within a 2D system, winding would build up in the transverse direction (cf. phase slips in 2D superfluids)

# Vorticity-winding transmutation



Net winding:

$$Q = \oint_{\partial\Omega} \frac{d\theta}{2\pi} \quad \mathbb{Z}^2$$

Stokes

$$\rightarrow \int_{\Omega} d^2r \rho_v$$

where

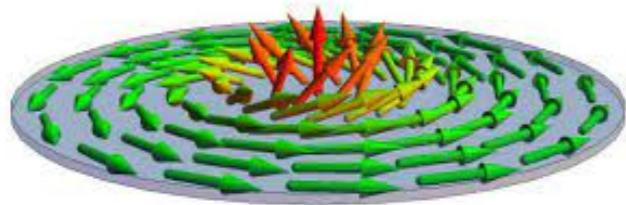
$$\rho_v = \frac{\mathbf{e}_z \cdot \partial_x \mathbf{n} \times \partial_y \mathbf{n}}{A}$$

vorticity density

$$\mathbf{j}_v^{(z)} = \frac{\epsilon^{ij} \mathbf{e}_z \cdot \partial_i \mathbf{n} \times \partial_j \mathbf{n}}{A}$$

vorticity flux

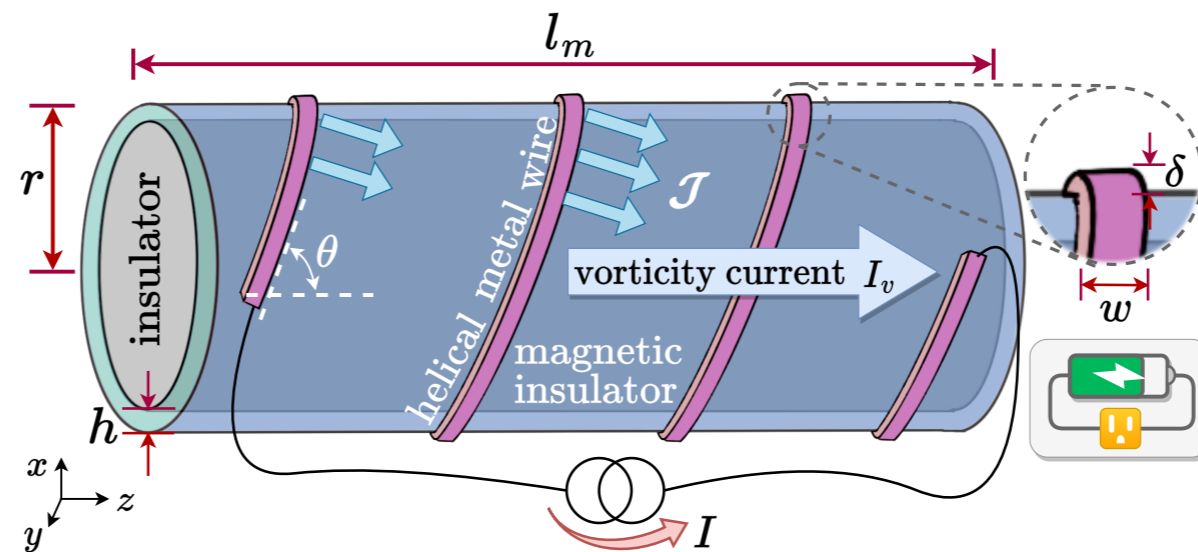
net integrated vorticity is naturally conserved



$$Q = \int d^2r \rho_v \rightarrow 1$$

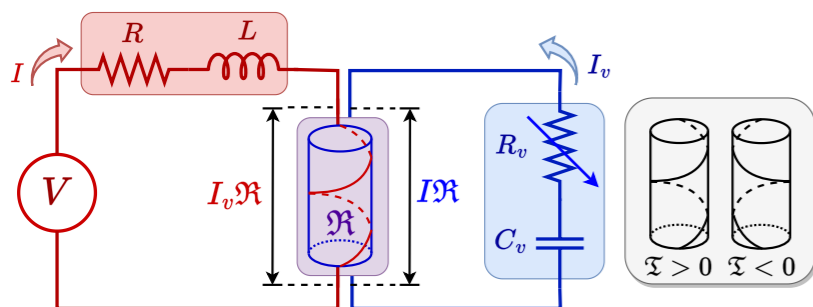
topological charge  $Q$  labels winding homotopy classes in the strongly-ordered XY limit

# Geometry-enabled topological energy storage



spin-transfer input power:  $\dot{W} = \int dl \boldsymbol{\tau} \cdot \mathbf{m} \times \partial_t \mathbf{m} = \mathfrak{R} I I_v$

In this geometry, electric current along the spiral generically drives vorticity flow along the cylinder axis (thus building up uniform transverse winding)



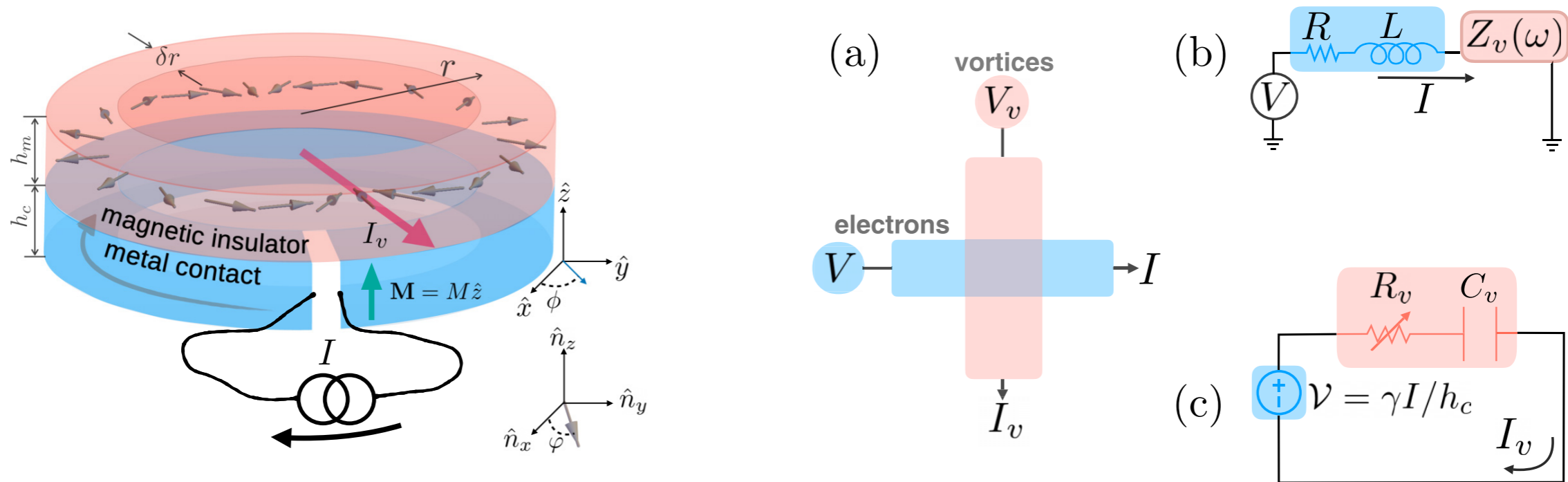
$$\xi = \mathfrak{R}^2 / R R_v$$

electron-vortex  
"cooperativity"

the effective dimensionless parameter, which is thermodynamically bounded to  $[0, 1]$ , is formally analogous to the thermoelectric figure of merit called  $ZT$



# Onsager description of the magnetic energy storage



- Vortex and electron fluxes cross-couple via Magnus-like friction:

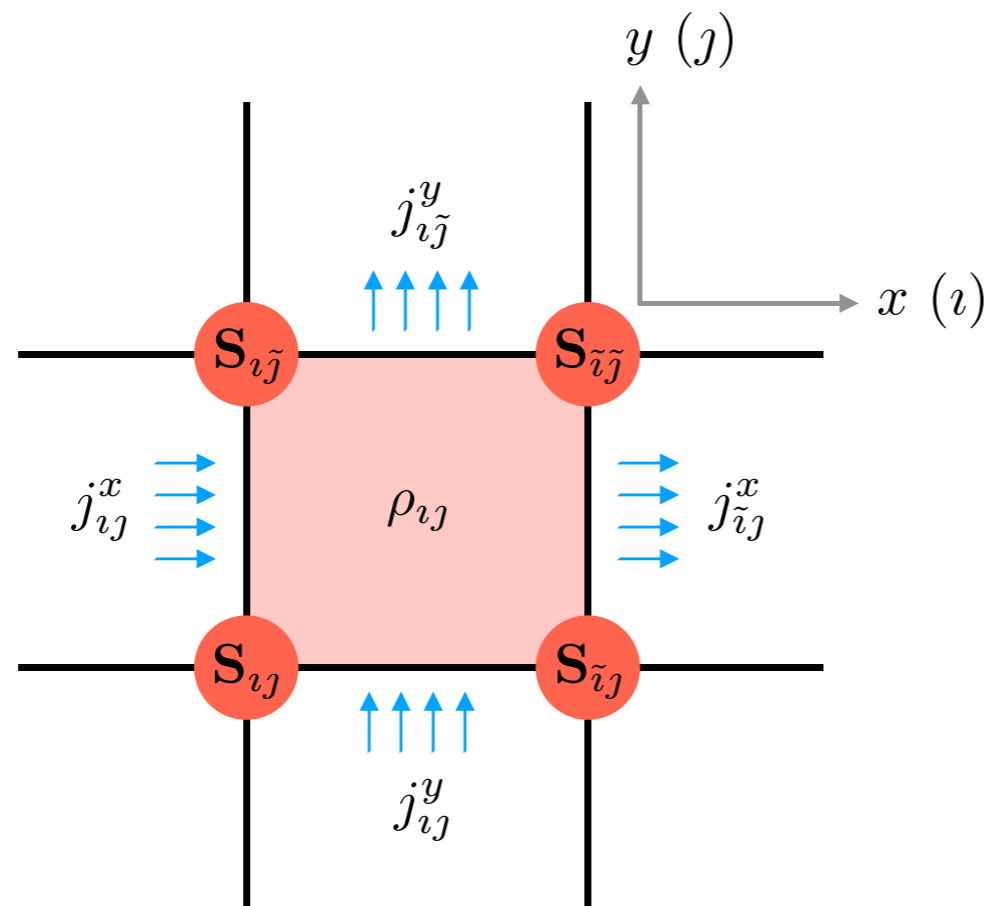
$$\begin{pmatrix} V \\ V_v \end{pmatrix} = \begin{pmatrix} R & \mathfrak{R} \\ -\mathfrak{R} & R_v \end{pmatrix} \begin{pmatrix} I \\ I_v \end{pmatrix}$$

- The magnetic annulus acts as a winding capacitor:

$$E_v = \frac{Q_w^2}{2C_v} \quad I_v = \frac{dQ_w}{dt}$$

- It is tempting to think of devices where a tunable vortex conductivity (e.g., associated with their binding and/or defect pinning) is utilized as a switch

# The vorticity is robust against quantum fluctuations



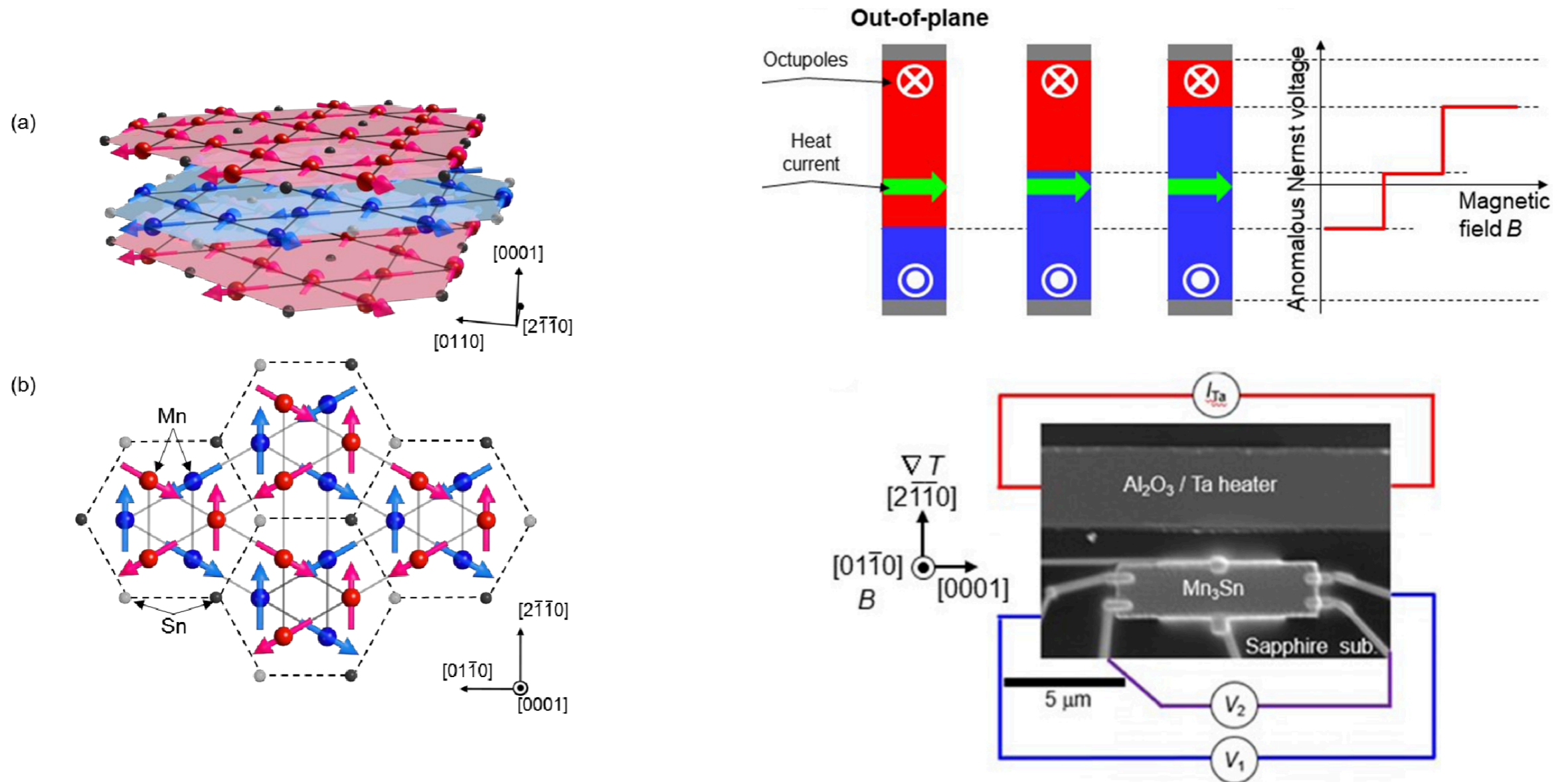
$$\partial_t \rho_{ij} + \frac{j_{i\tilde{j}}^x - j_{ij}^x + j_{ij}^y - j_{i\tilde{j}}^y}{a} = 0$$

- Vorticity per plaquette is given by the (z component of the) vector chirality: quantum version of the winding around the plaquette

$$\rho_{ij} = \frac{\mathbf{z} \cdot \mathbf{c}_{ij}}{2\pi a^2} \quad \text{where} \quad \mathbf{c}_{ij} \equiv \frac{1}{S^2} \sum_l \mathbf{S}_l \times \mathbf{S}_{\tilde{l}}$$

- For smooth classical textures, this reproduces the previous continuum version

# A new promising materials platform: $Mn_3Sn(Ge)$



Antiferromagnetic Weyl semimetal with planar (octupolar) spin texture free to rotate within the easy plane

Liu and Balents, *PRL* (2017)

Nernst measurement of a domain-wall motion subjected to a magnetic field

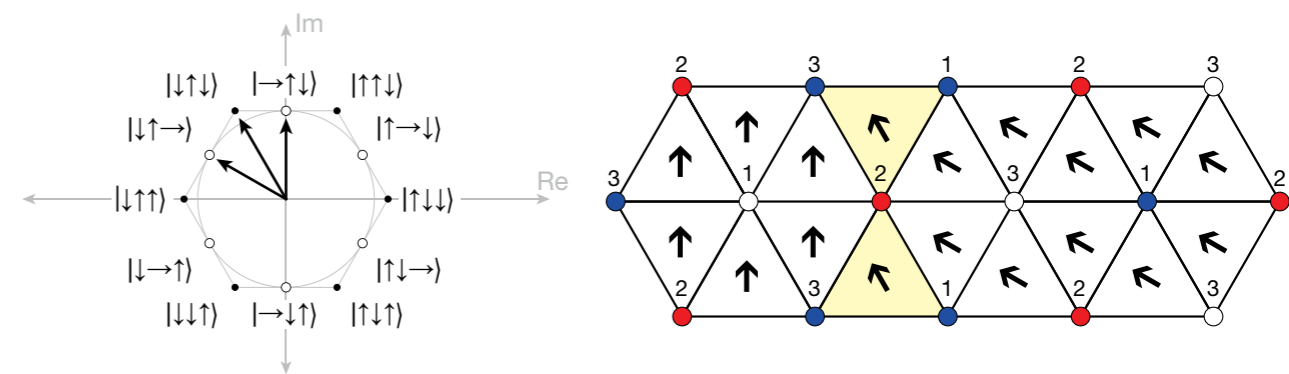
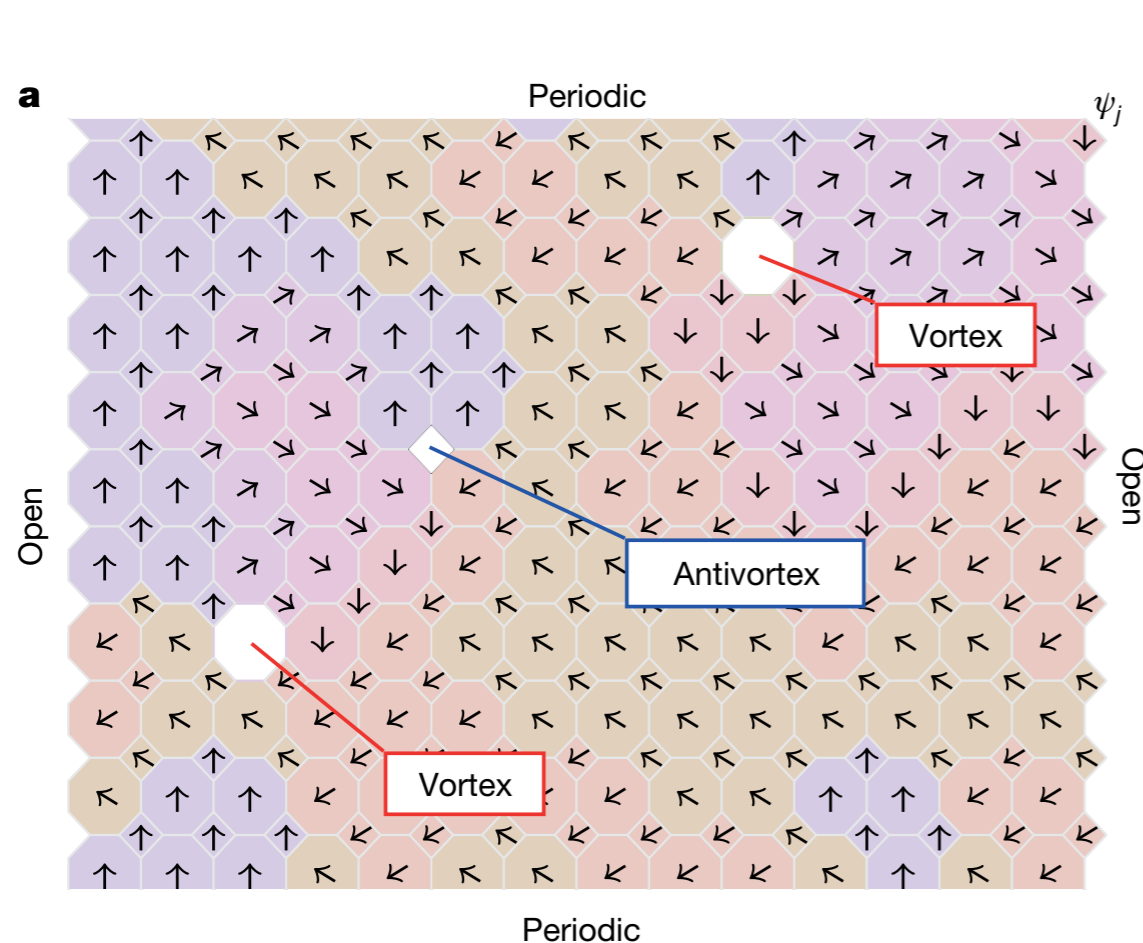
Otani and Higo, *APL* (2021)

# Quantum annealing of the Kosterlitz-Thouless transition

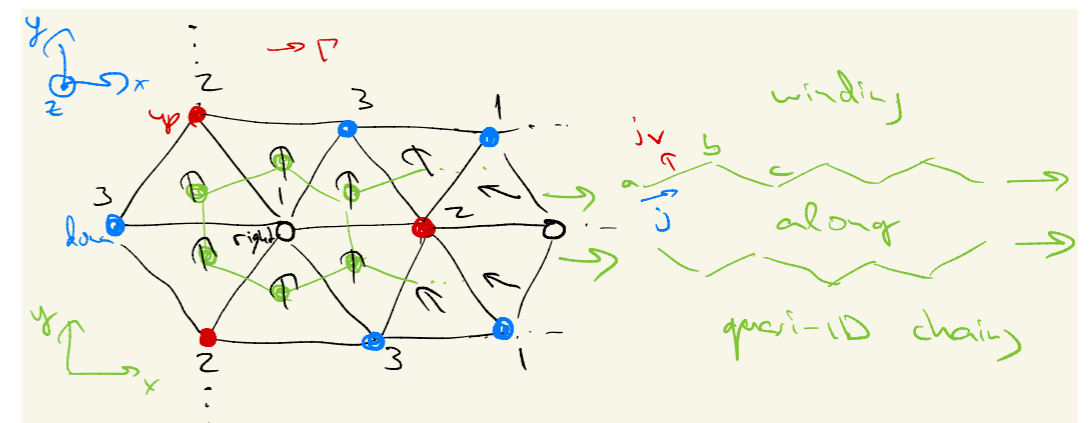
## Observation of topological phenomena in a programmable lattice of 1,800 qubits

Andrew D. King<sup>1\*</sup>, Juan Carrasquilla<sup>2</sup>, Jack Raymond<sup>1</sup>, Isil Ozfidan<sup>1</sup>, Evgeny Andriyash<sup>1</sup>, Andrew Berkley<sup>1</sup>, Mauricio Reis<sup>1</sup>, Trevor Lanting<sup>1</sup>, Richard Harris<sup>1</sup>, Fabio Altomare<sup>1</sup>, Kelly Boothby<sup>1</sup>, Paul I. Bunyk<sup>1</sup>, Colin Enderud<sup>1</sup>, Alexandre Fréchet<sup>1</sup>, Emile Hoskinson<sup>1</sup>, Nicolas Ladizinsky<sup>1</sup>, Travis Oh<sup>1</sup>, Gabriel Poulin-Lamarre<sup>1</sup>, Christopher Rich<sup>1</sup>, Yuki Sato<sup>1</sup>, Anatoly Yu. Smirnov<sup>1</sup>, Loren J. Swenson<sup>1</sup>, Mark H. Volkmann<sup>1</sup>, Jed Whittaker<sup>1</sup>, Jason Yao<sup>1</sup>, Eric Ladizinsky<sup>1</sup>, Mark W. Johnson<sup>1</sup>, Jeremy Hilton<sup>1</sup> & Mohammad H. Amin<sup>1,3</sup>

456 | NATURE | VOL 560 | 23 AUGUST 2018



$$\psi = \hat{\sigma}_1^z + \hat{\sigma}_2^z e^{2\pi i/3} + \hat{\sigma}_3^z e^{-2\pi i/3}$$



$$T < T_{KT}$$

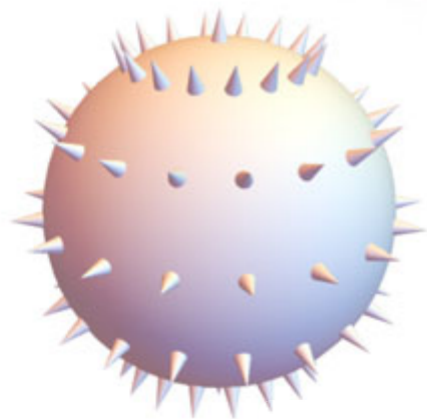
vortices bind with antivortices:  
vorticity insulator

$$T > T_{KT}$$

(anti)vortices unbind and proliferate:  
vorticity metal (two-component plasma)

# 3D version of vortices: Hedgehogs

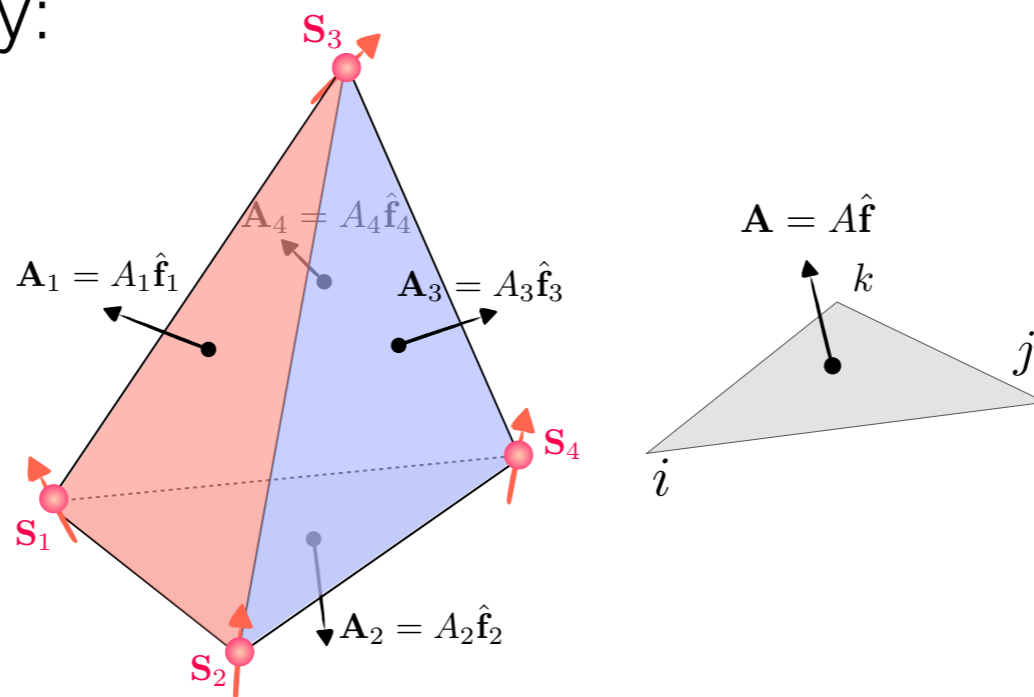
- The continuum hedgehog hydrodynamics



$$j^\mu = \epsilon^{\mu\alpha\beta\gamma} \partial_\alpha \mathbf{n} \cdot (\partial_\beta \mathbf{n} \times \partial_\gamma \mathbf{n}) / 8\pi, \quad \partial_\mu j^\mu = 0$$

$$\mathbf{n}(\mathbf{r}, t) : \quad \mathbf{n}, \mathbf{r} \in \mathbb{R}^3$$

has a natural quantum underpinning on an arbitrary lattice, in terms of scalar spin chirality:

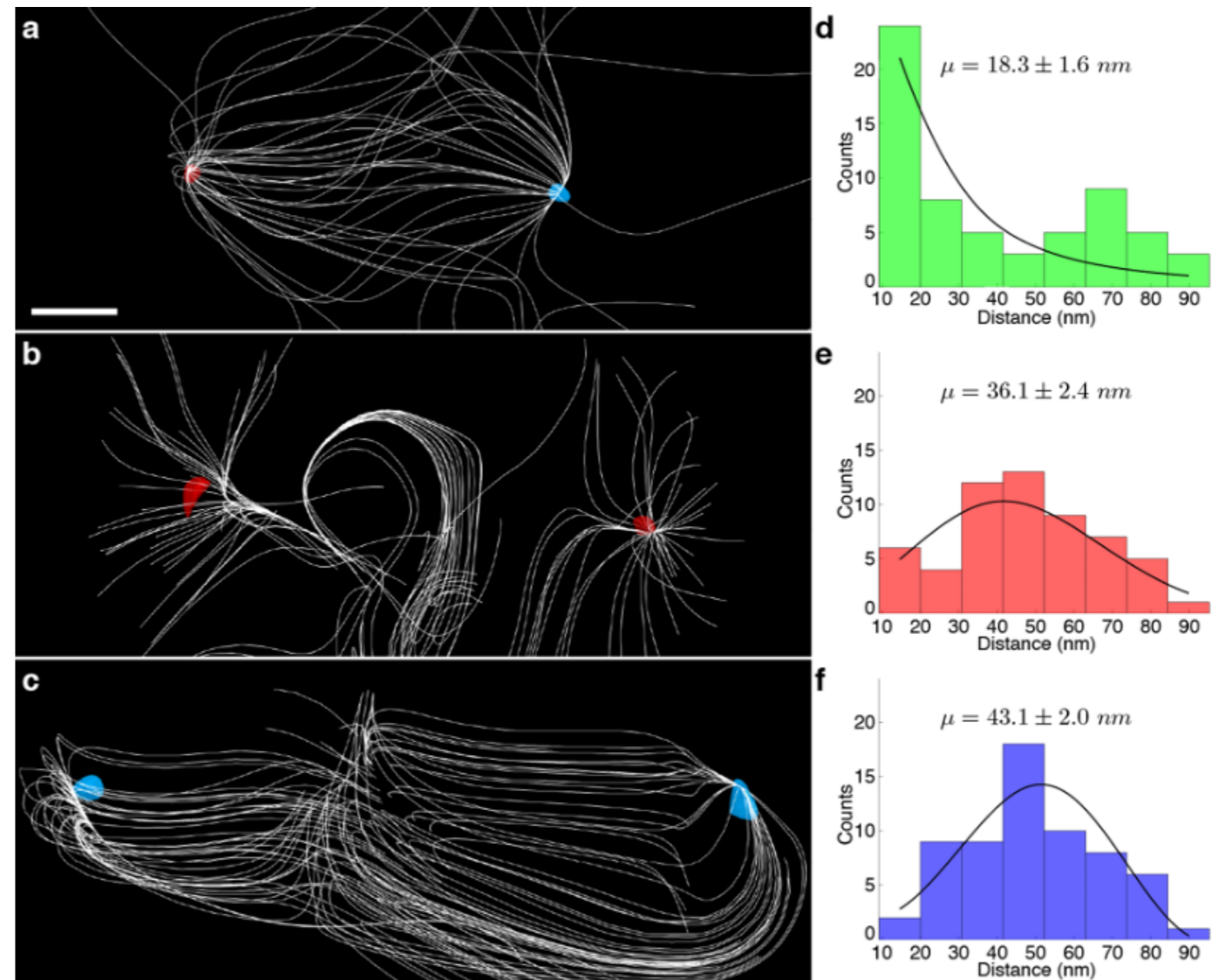
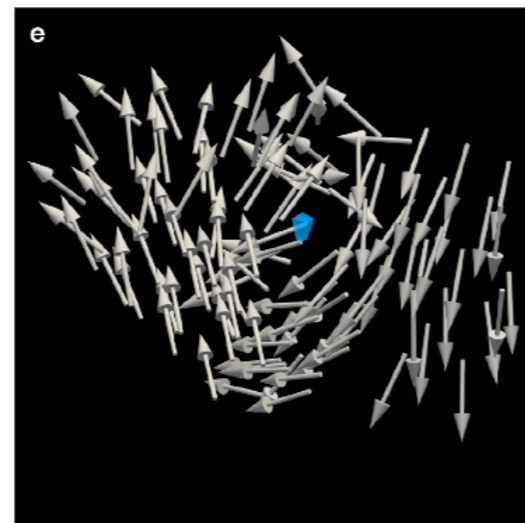
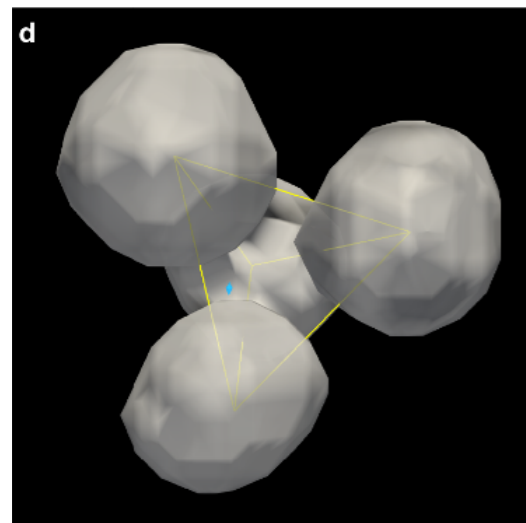
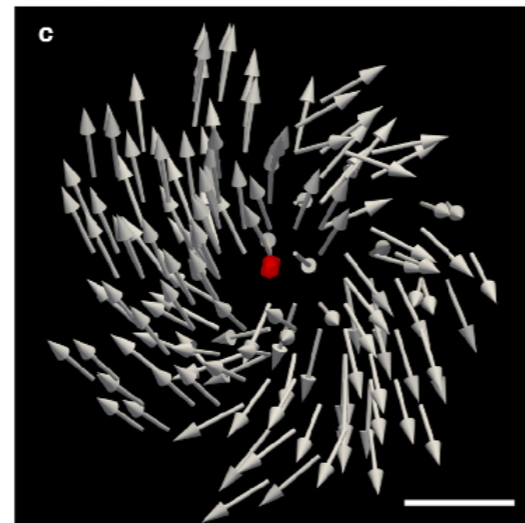
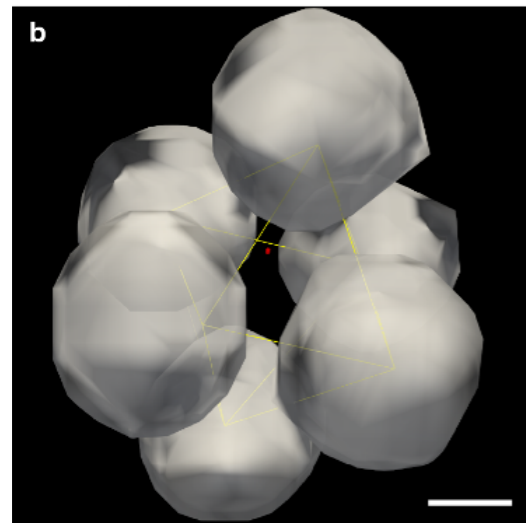
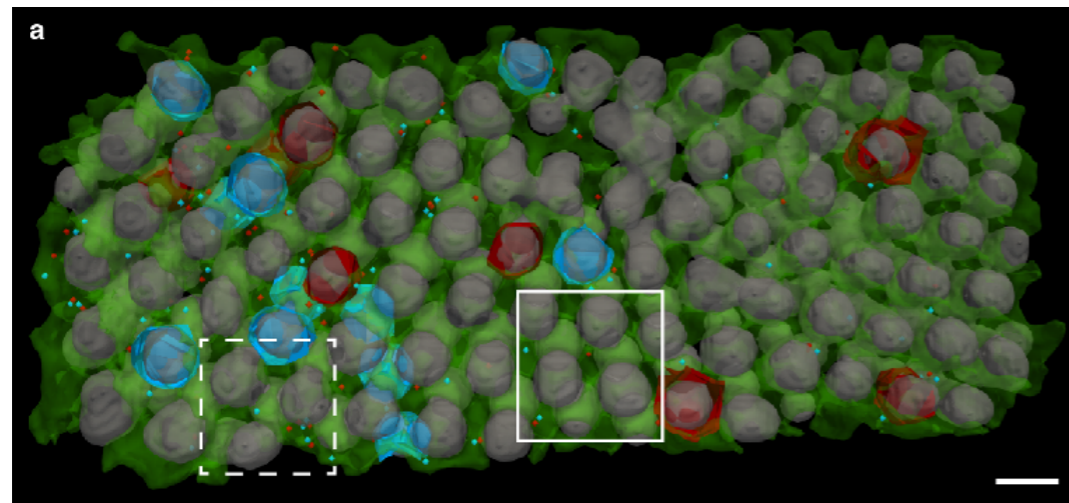


$$\rho = \frac{c_{123} + c_{142} + c_{243} + c_{134}}{8\pi V}$$

where

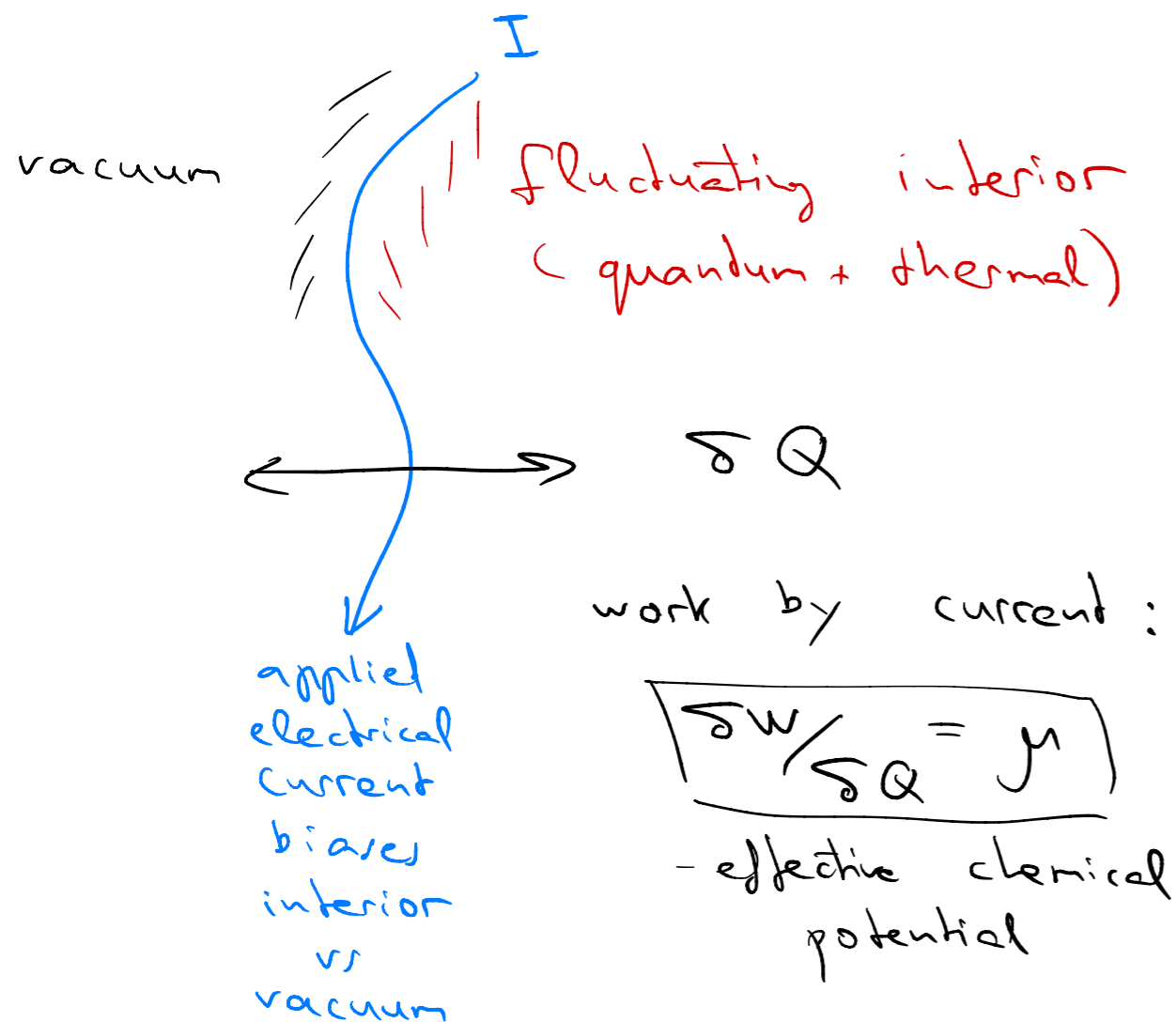
$$c_{ijk} = \frac{\mathbf{S}_i \cdot \mathbf{S}_j \times \mathbf{S}_k}{S^3}$$

# Engineering and imaging hedgehog metamaterials



Creation and observation of topological magnetic monopoles and their interactions in a ferromagnetic metal lattice  
Soft X-ray imaging by Miao's group at UCLA [Rana et al., Nature Nano. (2023)]

# Electrical injection (detection) of topological charge $\delta Q$



- A nonequilibrium flow of  $Q$  must be consistent with crystalline and structural symmetries, Onsager reciprocity, and (1st/2nd) laws of thermodynamics
- In principle, we generally expect such injection to happen for a generic drive (electrical/microwave/optical etc.), as long as it is not disallowed by symmetries

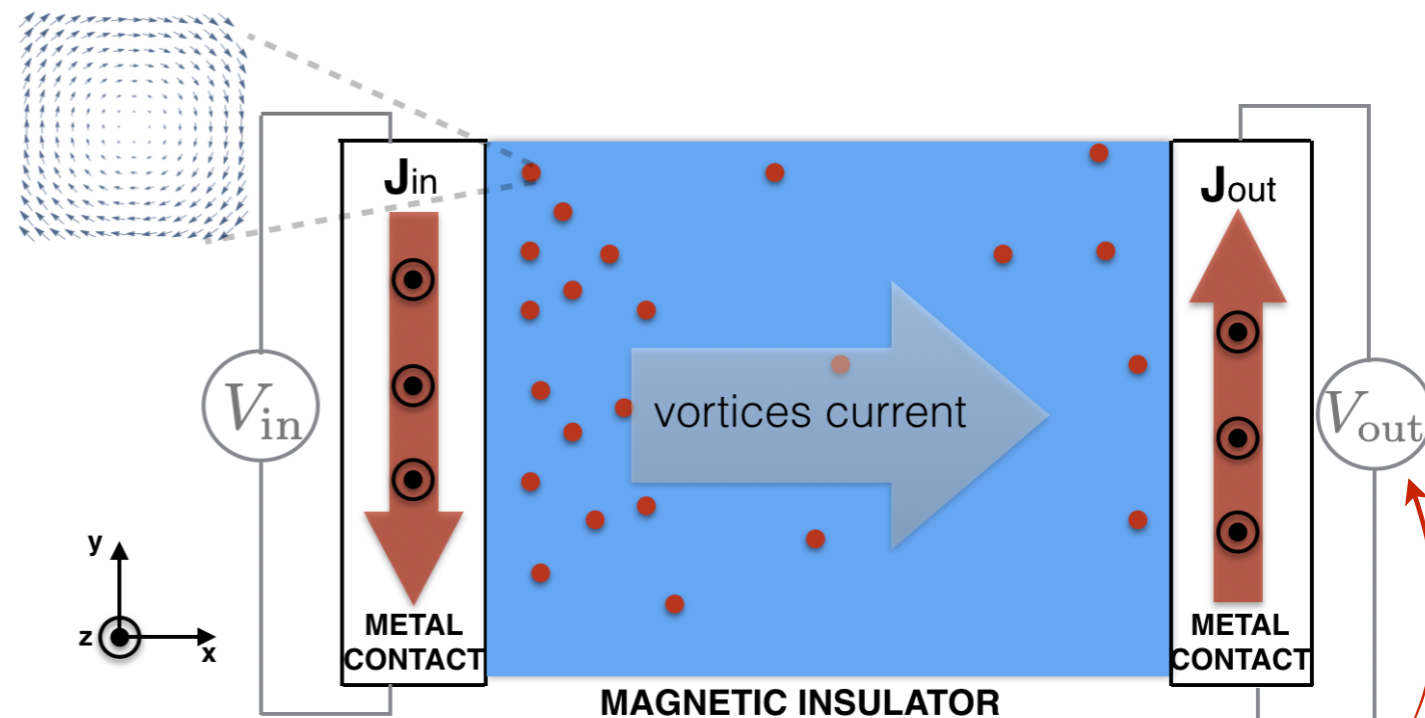
Pierre Curie: Asymmetry creates the phenomenon

# Electrical injection (detection) of vorticity

Energetics (thermodynamics) and symmetries:

spin torque by the input current at the left interface:

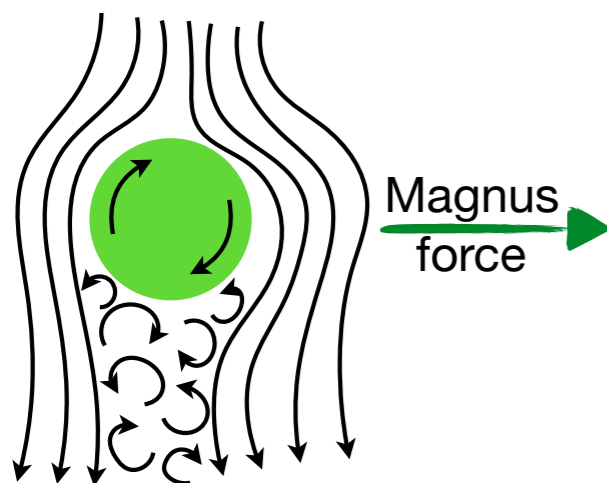
$$\boldsymbol{\tau} = \eta \mathbf{M} \cdot \mathbf{m} (\mathbf{J} \cdot \nabla) \mathbf{m}$$



$\mathbf{M} \propto \mathbf{z}$  - magnetization of the metal contact

$\mathbf{m}(x, y, t)$  - magnetic texture of the insulator

Onsager-reciprocal motive force



$$\dot{W} = \int dy \boldsymbol{\tau} \cdot (\mathbf{m} \times \partial_t \mathbf{m}) \rightarrow \eta \mathbf{z} \cdot \mathbf{J} \times \mathbf{j}$$

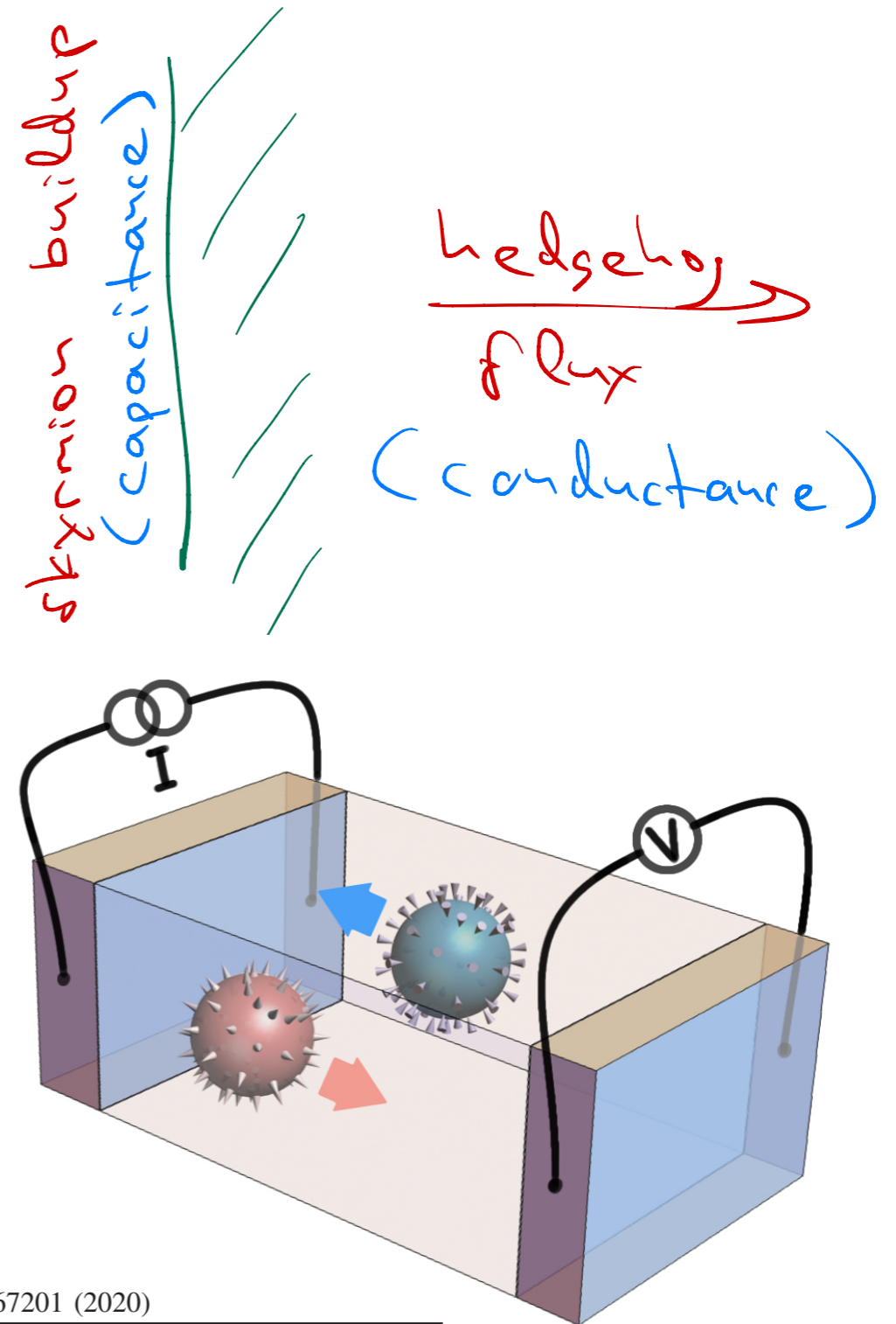
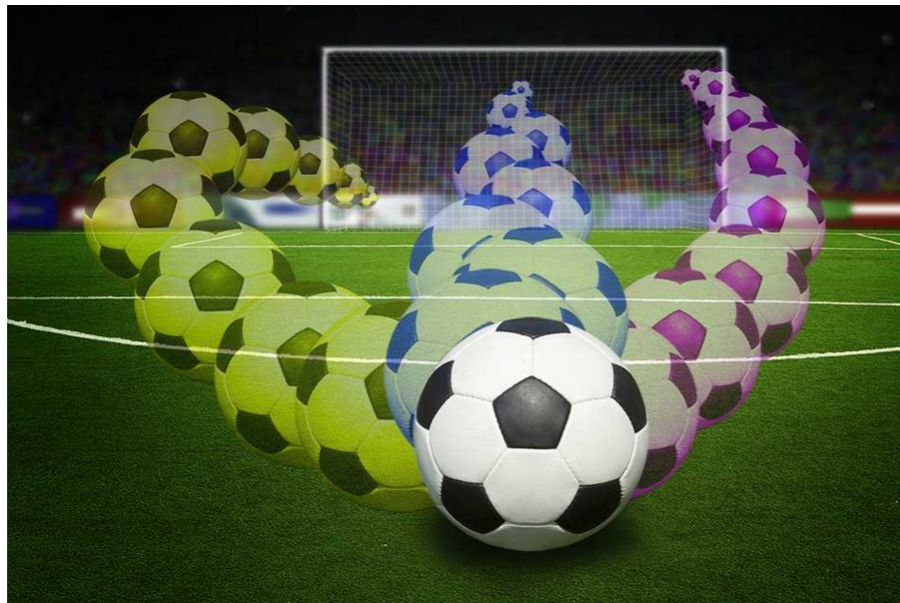
charge current

vortex flux



# Extending to 3D bulk: Hedgehog flow

The 3D hedgehog flow with skyrmionic bulk/boundary correspondence is closely analogous to the 2D vorticity flow with winding bulk/boundary correspondence



PHYSICAL REVIEW LETTERS **125**, 267201 (2020)

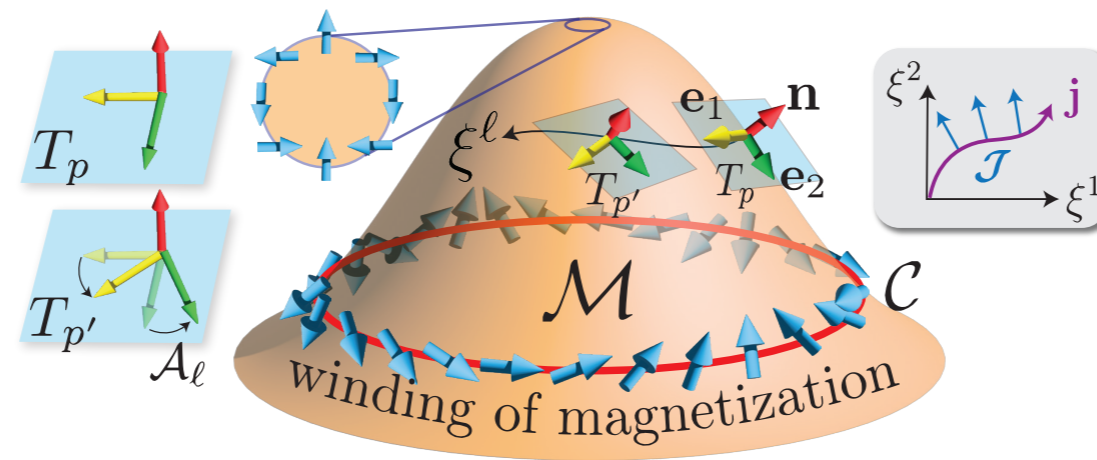
Editors' Suggestion

## Topological Transport of Deconfined Hedgehogs in Magnets

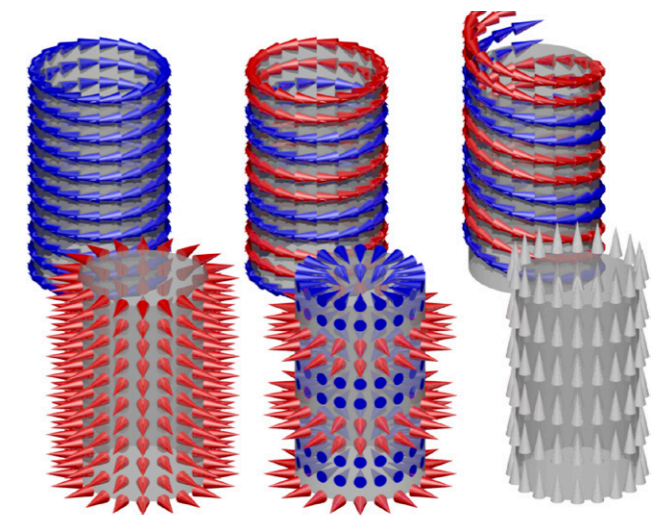
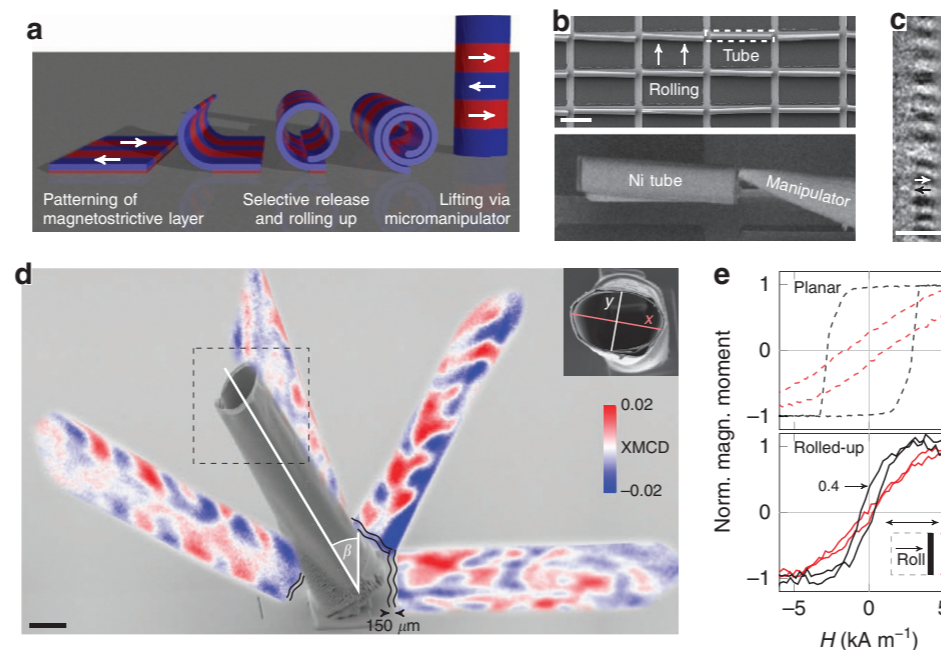
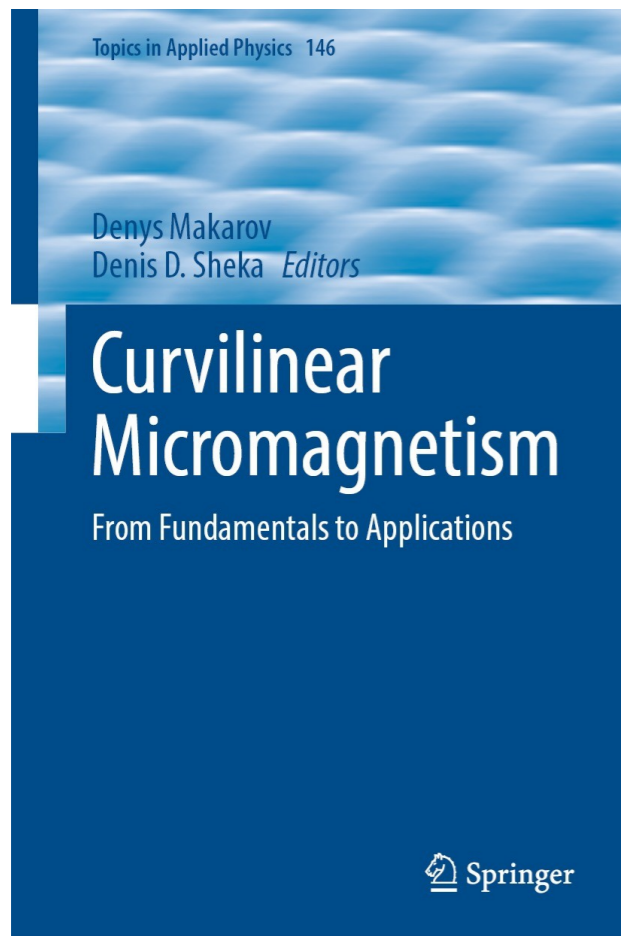
Ji Zou<sup>✉</sup>, Shu Zhang,<sup>\*</sup> and Yaroslav Tserkovnyak

Department of Physics and Astronomy, University of California, Los Angeles, California 90095, USA

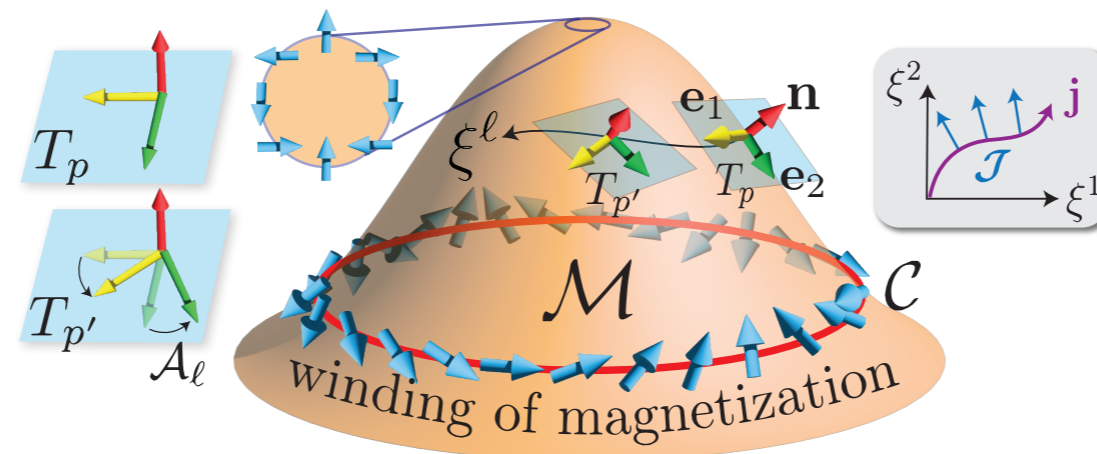
# Transport of vorticity on curved surfaces



- Can nontrivial geometry reduce the symmetries enough to allow to drive topological hydrodynamics even in a single magnetic layer?

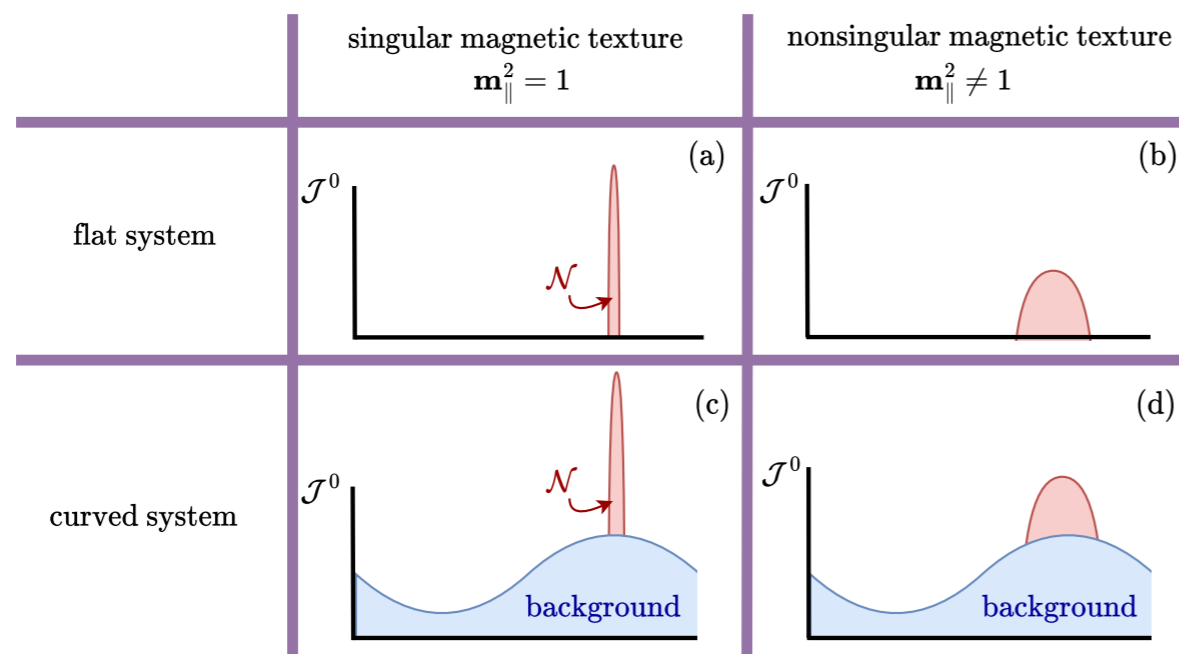


# Hydrodynamics on curved surfaces



- Net topological charge is defined in terms of a winding 1-form (the density is constructed via exterior derivative, utilizing generalized Stokes' theorem):

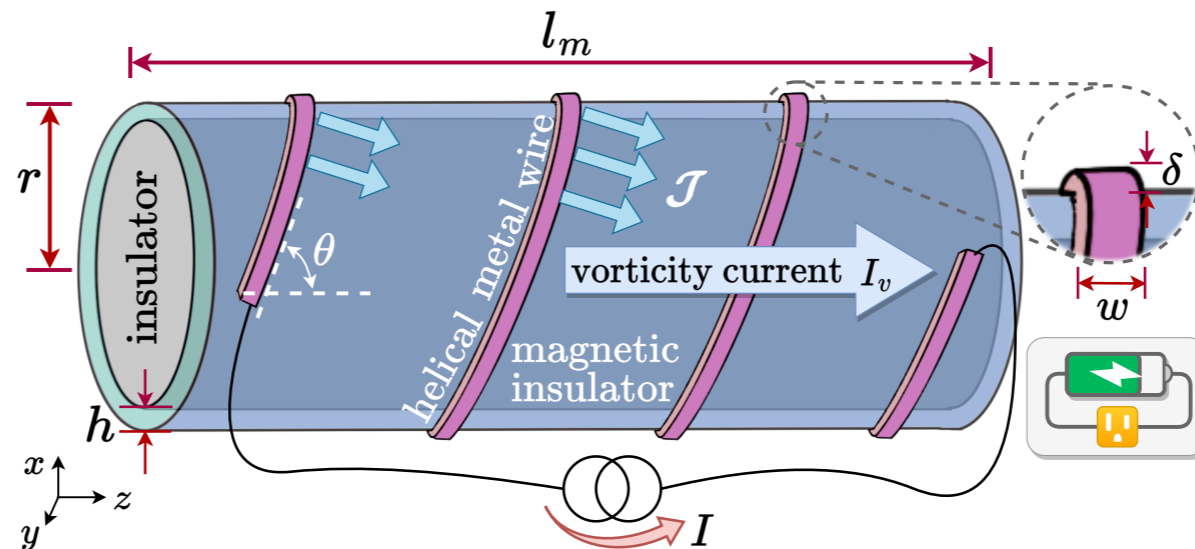
$$Q = \int_{\partial S} \rho_w = \frac{1}{2\pi} \int_{\partial S} d\xi^i \mathbf{m}_{\parallel}^2 D_i \varphi \quad \text{where} \quad D_i \varphi \equiv \partial_i \varphi - \mathbf{e}_1 \cdot \partial_i \mathbf{e}_2$$



in the strongly easy-plane limit:

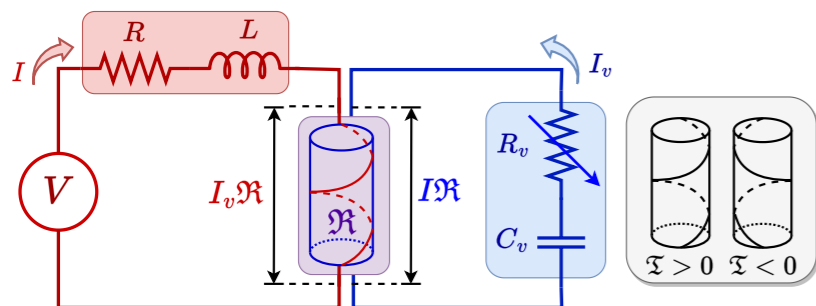
$$Q = \mathcal{N} - \frac{1}{2\pi} \int_S d\xi^1 d\xi^2 \sqrt{g} \mathcal{K} \quad (\text{Mermin-Ho})$$

# “Simplified” topological energy storage



spin-transfer input power:  $\dot{W} = \int dl \boldsymbol{\tau} \cdot \mathbf{m} \times \partial_t \mathbf{m} \propto \mathfrak{T} I I_v$

$\mathfrak{T}$  - (pseudoscalar) “torsion of a curve” (electrical wire), which can effectively (from the symmetry point of view) convert the local geometric normal to the surface into an out-of-plane magnetization



$$\xi \propto \mathfrak{T}^2 / R R_v$$

the effective dimensionless parameter, which is thermodynamically bounded to  $[0, 1]$ , is formally analogous to the thermoelectric figure of merit called  $ZT$

# Outlook

---

- ◆ Dynamics of collective order-parameter textures can have robust low-energy behavior rooted in topological conservation laws and responsive to geometric controls
- ◆ Spin-based systems are abundant, versatile, and amenable to the wealth of spintronic tools
- ◆ This can lead to new strategies for probing materials as well as applications, such as energy storage and nontraditional computing (both classical and quantum)
- ◆ Myriad connections across different fields of physics, from astrophysics to turbulence
- ◆ On the quantum front, intriguing outlooks concern direct transport probes of condensed-matter dualities (e.g., vortex condensation at the superfluid-insulator transition), interplay between real-space and momentum-space topologies, and integration with optically-active quantum impurities for sensing and generation of quantum entanglement

